

Non-compact Spin(7)-manifolds

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Spin(7)-manifolds

- $\text{Spin}(7) = \text{Stab}_{\text{GL}(8, \mathbb{R})}(\psi_0)$, where the model 4-form $\psi_0 \in \Lambda^4(\mathbb{R}^8)^*$ is given by

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- $G_2 \subset \text{Spin}(7)$: If we write $\mathbb{R}^8 = \mathbb{R}_t \times \mathbb{R}^7$ and $\varphi_0 \in \Lambda^3(\mathbb{R}^7)^*$ is the standard G_2 -structure on \mathbb{R}^7 , then

$$\psi_0 = dt \wedge \varphi_0 + *7\varphi_0.$$

- A 4-form $\psi \in \Omega(M)$ on an 8-manifold M is called *admissible* if at each point $p \in M$ there is a linear isomorphism $T_p M \cong \mathbb{R}^8$ which identifies $\psi|_p = \psi_0$

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- Restrictions on the holonomy group put algebraic constraints on the curvature tensor. $\text{Hol}(g) \subset \text{Spin}(7)$ implies $\text{Ric}(g) = 0$.

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- $d\psi = 0$ locally non-linear PDE system of 56 equations for 43 unknown functions depending on 8 variables.
- Idea: Use symmetries to simplify.
- There are some restrictions: There are no non-flat homogeneous Spin(7)-manifolds and no compact Spin(7)-holonomy manifolds with continuous symmetries.

Cohomogeneity one manifolds

- Let G be a compact Lie group. G acts on M with *cohomogeneity one* if M/G is 1-dimensional.

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- In the context of $\text{Spin}(7)$ -manifolds only one topological type of cohomogeneity one manifolds is interesting:
- $M/G \cong [0, \infty)$
- Denote the quotient map by $q : M \rightarrow M/G$. Then

$$M = \underbrace{(G/K)}_{q^{-1}(0)} \cup \underbrace{(0, \infty) \times (G/H)}_{q^{-1}(0, \infty)}.$$

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- G/K is called *singular orbit* and G/H is called *principal orbit*.

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- Simplest example: $\mathbb{R}^n = \{0\} \cup (0, \infty) \times S^{n-1}$, where $S^{n-1} = \text{SO}(n)/\text{SO}(n-1)$.
- This is the local model for the general case:
total space: vector space \rightarrow vector bundle
principal orbits: spheres \rightarrow sphere bundles
singular orbit: zero \rightarrow zero section

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- Suppose M is a 8-dimensional cohomogeneity one manifold as above and ψ a G -invariant Spin(7)-structure.
- On each principal orbit $\{t\} \times G/H$ there exists a G -invariant G_2 -structure φ_t such that

$$\begin{aligned}\psi &= dt \wedge \varphi_t + *_t \varphi_t, \\ g &= dt^2 + h_t.\end{aligned}$$

Here the Hodge star depends on φ_t . The condition $d\psi = 0$ for ψ to be torsion-free then is equivalent to the system

$$\begin{aligned}d_{G/H} * \varphi_t &= 0, \\ \frac{\partial}{\partial t} * \varphi_t &= d_{G/H} \varphi_t.\end{aligned}$$

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- To get examples we can consider initial value problems: if we prescribe a coclosed G -invariant G_2 -structure on $\{t_0\} \times G/H$ then there exists a unique solution on $(t_0 - \varepsilon, t_0 + \varepsilon) \times G/H$.
- Eschenburg–Wang (2000) in the context of cohom1 Einstein manifolds considered a singular initial value problem on the singular orbit. In general one obtains families of local solution in a neighbourhood of the singular orbit which depend on several free parameters.

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- $\Psi_\mu, \mu \in (0, \infty)$, has principal orbit $N(1, -1) \cong \text{SU}(3)/\text{U}(1)$ and singular orbit $S^5 = \text{SU}(3)/\text{SU}(2)$. Total space is unique non-trivial rank 3 vector bundle over S^5 .

- Reidegeld ('08) has carried out this programme in the context of cohomological Spin(7)-manifolds. Among other things he finds two 1-parameter (up to scale) families:
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- $\Upsilon_\tau, \tau \in (-\infty, \infty)$, has principal orbit $N(1, 0) \cong \mathrm{SU}(3)/\mathrm{U}(1)$ and singular orbit $\mathbb{C}P^2 \cong \mathrm{SU}(3)/(\mathrm{SU}(2) \times \mathrm{U}(1))$. Total space is universal quotient bundle of $\mathbb{C}P^2$.

- For every pair (k, l) of integers which are not both zero, $U(1)$ can be embedded in the maximal torus of diagonal matrices in $SU(3)$ as

$$e^{i\theta} \mapsto \begin{pmatrix} e^{ik\theta} & 0 & 0 \\ 0 & e^{il\theta} & 0 \\ 0 & 0 & e^{-i(k+l)\theta} \end{pmatrix}. \quad (1.1)$$

We also denote this subgroup of $SU(3)$ by $U(1)_{k,l}$. The Aloff–Wallach space $N(k, l)$ is the homogeneous space $SU(3)/U(1)_{k,l}$.

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- An important property of complete non-compact manifolds is the asymptotic geometry.

Asymptotic Geometry

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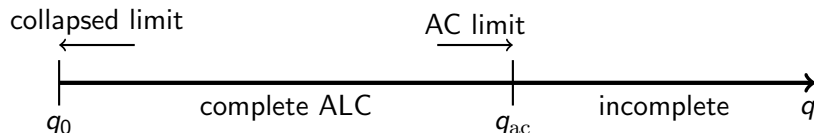
- $(C(\Sigma), \psi)$ is a Spin(7)-cone iff (Σ, φ) is a *nearly parallel* G_2 -manifold.
- Every Aloff–Wallach space $N(k, l)$ carries a homogeneous, *nearly parallel* G_2 -structure.

- The other asymptotic type that we will meet are *asymptotically locally conical* $\text{Spin}(7)$ -manifolds. Close to infinity they locally look like the product of a 7-dimensional cone and a circle of fixed length.
- These are analogues of *asymptotically locally flat* (ALF) hyperkähler manifolds in real dimension 4 such as the Taub-NUT and Atiyah–Hitchin metrics.

Behaviour of the families Ψ_μ and Υ_τ

Theorem (L-, '20)

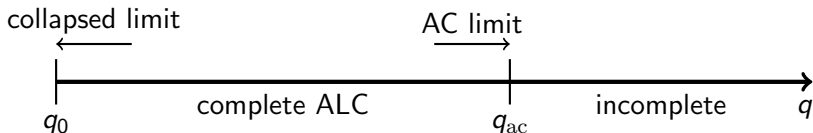
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As $\mu \rightarrow 0$ the curvatures stay bounded and Ψ_μ converges to the Bryant–Salamon G_2 -holonomy metric on $\Lambda_-^2 \mathbb{C}P^2$.

As $\tau \rightarrow -\infty$ the curvatures of Υ_τ blow up on the zero section.

- The AC Spin(7) holonomy manifolds $\Psi_{\mu_{ac}}$ and $\Upsilon_{\tau_{ac}}$ resemble the well-known conifold transition of Calabi-Yau 3-folds: the *smoothing* T^*S^3 of the conifold $\{(z_1, z_2, z_3, z_4) \in \mathbb{C}^4 \mid z_1^2 + z_2^2 + z_3^2 + z_4^2 = 0\}$ and its *small resolution* $\mathcal{O}(-1) \oplus \mathcal{O}(-1) \rightarrow \mathbb{C}P^1$ are topologically different spaces which carry cohomogeneity one AC Calabi-Yau metrics asymptotic to the same Calabi-Yau cone, the conifold.

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- The conifold transition was predicted by Gukov–Sparks–Tong ('03)
- For each τ the singular orbit $\mathbb{C}P^2$ is Cayley with respect to Υ_τ .

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- $N(1, -1) \cong SU(3)/U(1)$ is a circle bundle over the flag manifold $SU(3)/U(1)^2$. f describes the size of the fibres.
- (a, b, c) describe a metric on $SU(3)/U(1)^2$.
- The Bryant–Salamon metric on $\Lambda_-^2 \mathbb{C}P^2$ is asymptotic to the cone over the nearly Kähler metric $(a, b, c) = (1, 1, 1)$ on the flag manifold. Therefore, ALC asymptotics correspond to

$$(a(t), b(t), c(t), f(t)) \rightarrow (t, t, t, \ell) \quad \text{as } t \rightarrow \infty,$$

where $\ell > 0$ is a constant.

- Denote by (a_c, b_c, c_c, f_c) the coefficients of the nearly G_2 -metric on $N(1, -1)$. Then AC asymptotics correspond to

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- $d\psi = 0$ then gives the ODE system

$$\begin{aligned} \frac{\dot{a}}{a} &= \frac{b^2 + c^2 - a^2}{abc}, \\ \frac{\dot{b}}{b} &= \frac{c^2 + a^2 - b^2}{abc} - \frac{f}{b^2}, \\ \frac{\dot{c}}{c} &= \frac{a^2 + b^2 - c^2}{abc} + \frac{f}{c^2}, \\ \frac{\dot{f}}{f} &= \frac{f}{b^2} - \frac{f}{c^2}. \end{aligned}$$

Projective coordinates

- Idea: The right-hand side of the ODE system is a scale-invariant rational expression. Switch to projective coordinates

$$A = \frac{a}{c}, \quad B = \frac{b}{c}, \quad F = \frac{f}{c}.$$

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$$A = \frac{a}{c}, \quad B = \frac{b}{c}, \quad F = \frac{f}{c}.$$

- With respect to the parameter s given by $dt = \frac{ab}{c} ds$ the ODE system for (A, B, F) takes the form

$$\begin{aligned}\dot{A} &= A(2 - 2A^2 - ABF), \\ \dot{B} &= B\left(2 - 2B^2 - ABF - \frac{AF}{B}\right), \\ \dot{F} &= F\left(1 - A^2 - B^2 - 2ABF + \frac{AF}{B}\right).\end{aligned}$$

Good news:

- If we can solve the system for (A, B, F) , we can always solve the final remaining ODE to obtain a solution (a, b, c, f) .

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- The asymptotic geometries now correspond to critical points.
- The ODE system for (A, B, F) is less singular.

Bad news:

- Still no control on F .

Another coordinate transform

- Key: Well-behaved behaviour of A, B dominates bad-behaviour of F .

Another coordinate transform

- Key: Well-behaved behaviour of A, B dominates bad-behaviour of F .
- A good set of coordinates to approach the problem is

$$X = A^2, \quad Y = B^2, \quad Z = ABF.$$

Still denoting differentiation with respect to the variable s by a dot, the ODE system takes the form

$$\begin{aligned}\dot{X} &= 2X(2 - 2X - Z), \\ \dot{Y} &= 4Y - 4Y^2 - 2YZ - 2Z, \\ \dot{Z} &= Z(5 - 3X - 3Y - 4Z).\end{aligned}$$

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- The family Ψ_μ of trajectories emanates from fixed point $(0, 1, 0)$ and enters \mathcal{W} .
- The family Υ_τ emanates from the fixed point $(1, 0, 0)$ and enters \mathcal{W} .

The fixed points in the closure of \mathcal{W} are:

- $(0, 1, 0)$: S^5 , index is $(2, 1)$
- $(1, 0, 0)$: $\mathbb{C}P^2$, index is $(2, 1)$
- $(1, 1, 0)$: ALC end, index is $(0, 3)$
- The $\text{Spin}(7)$ -cone

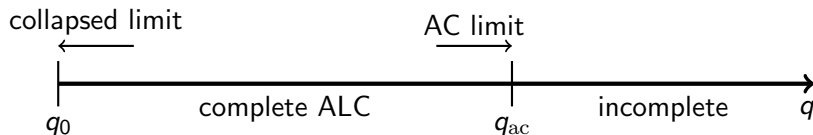
$$(X_c, Y_c, Z_c) := \left(\frac{15 - 3\sqrt{5}}{10}, \frac{3 - \sqrt{5}}{2}, \frac{3\sqrt{5} - 5}{5} \right) \approx (0.83, 0.38, 0.34)$$

index is $(1, 2)$.

- $(0, 0, 0)$: index is $(3, 0)$

Proof strategy

- Recall: We have the two families $\Psi_\mu, \mu \in (0, \infty)$ and $\Upsilon_\tau, \tau \in (-\infty, \infty)$ and want to show



- Problem: AC space is elusive.
- Define

$$\mathfrak{X}_{\text{alc}} := \{\mu \mid \Psi_\mu \text{ is complete and ALC}\},$$

$$\mathfrak{X}_{\text{ac}} := \{\mu \mid \Psi_\mu \text{ is complete and AC}\},$$

$$\mathfrak{X}_{\text{inc}} := \{\mu \mid \Psi_\mu \text{ is incomplete}\}.$$

Proof strategy: continuity principle

Our strategy is to prove

- $\mathfrak{X}_{\text{alc}}$ and $\mathfrak{X}_{\text{inc}}$ are both open and non-empty,
- $(0, \infty) = \mathfrak{X}_{\text{alc}} \cup \mathfrak{X}_{\text{ac}} \cup \mathfrak{X}_{\text{inc}}$.

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This then implies that \mathfrak{X}_a is non-empty! Moreover, a comparison argument then shows

$$\mathfrak{X}_{\text{alc}} = (0, \mu_1), \quad \mathfrak{X}_{\text{ac}} = [\mu_1, \mu_2], \quad \mathfrak{X}_{\text{inc}} = (\mu_2, \infty).$$

- ALC end $(1, 1, 0)$ is a sink $\Rightarrow \mathfrak{X}_{\text{alc}}$ is open

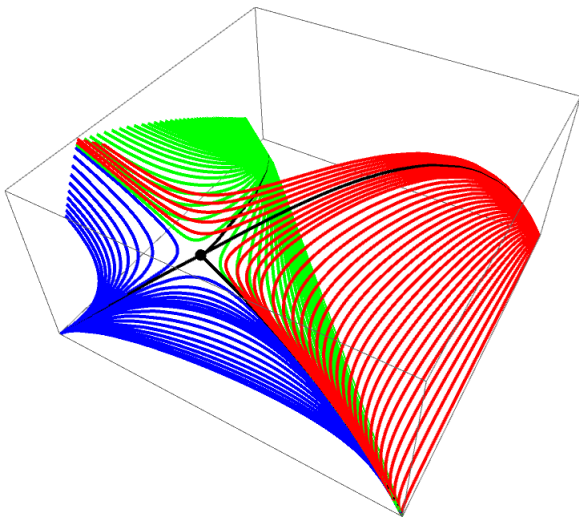
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- another key ingredient:
if a trajectory stays inside a bounded subset (such as \mathcal{W}), it must converge to a fixed point.
- There are only two fixed points in $\overline{\mathcal{W}}$ the trajectories Ψ_μ can converge: $\Rightarrow (0, \infty) = \mathfrak{X}_{\text{alc}} \cup \mathfrak{X}_{\text{ac}} \cup \mathfrak{X}_{\text{inc}}$



- Trajectory connecting $\text{Spin}(7)$ -cone and ALC end: conically singular ALC $\text{Spin}(7)$ holonomy metric

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- For $z \in (0, z_0)$ there exists 1-parameter family of Spin(7) holonomy metrics Ω_{κ}^z , $\kappa \in (0, 1)$, which behave as Ψ_{μ} and Υ_{τ} . Collapsed limit is singular version of Bryant–Salamon AC metric. While $a, b \rightarrow 0$ at singularity, $c \rightarrow \infty$.

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- Recent work by Chi on existence of ALC metrics

Let (M, ψ) AC Spin(7)-manifold with

$$|\psi - \psi_C| = \mathcal{O}(r^\nu).$$

ν is called decay rate.

- A tensor which decays faster than -4 is square integrable on (M, ψ) .
- Denote by \mathcal{M}_ν space of torsion-free AC Spin(7)-structures on M which decay at least like $r^{-\nu}$ to a fixed Spin(7)-cone, up to diffeomorphisms decaying to the identity.

Theorem (L-, '21)

In the non- L^2 regime ($-4 < \nu < 0$), for generic rates \mathcal{M}_ν is an orbifold. Dimension can be computed in terms of topology and solutions to differential equations on link.

- Derivative of

$$GL(7, \mathbb{R}) \rightarrow \Lambda^3(\mathbb{R}^7)^*, A \mapsto A^* \varphi_0$$

gives map

$$\mathfrak{gl}(7, \mathbb{R}) = \Lambda_7^2 \oplus \Lambda_{14}^2 \oplus \mathbb{R}Id \oplus \mathcal{S}_0^2(\mathbb{R}^7) \rightarrow \Lambda^3.$$

This gives decomposition

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- Infinitesimal deformations in the range $\nu \in (-4, 0)$ come from solutions of

$$d\zeta = -(\nu + 4) * \zeta, \quad \zeta \in \Gamma(\Lambda_{27}^3).$$

- Those forms also solve the eigenvalue equation

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- The nearly parallel G_2 -structures on the Aloff–Wallach spaces $N(k, l)$ are not standard. $SU(3)$ -invariance \rightarrow polynomial equation.