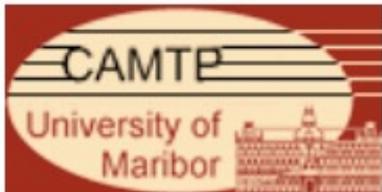


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Zero Modes of Higgs bundles on Spin(7) manifolds

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Motivation/Goals

- **Motivation:** M-theory/String Theory in four & three space-time dimensions with minimal (N=1) supersymmetry
[D=4 N=1 & D=3 N=1]
→ compactification on G_2 & Spin(7) holonomy spaces
- **Focus:** gauge degrees within G_2 & Spin(7) compactifications associated w/ co-dimension four (ADE) singularities, governed by Higgs bundles
- **Goal:** background solutions & zero modes (charged matter)
 - Higgs bundles on G_2 localized modes w/ gauge flux (“T-branes”)
[Barbosa, M.C., Heckman, Lawrie, Torres, Zoccarato 1906.02212]
 - Higgs bundles on Spin(7) - geometric unification w/ other bundles
[M.C., Heckman, Rochais, Torres, Zoccarato 2003.13682]
- zero modes: classical & quantum
[M.C., Heckman, Torres, Zoccarato 2107.00025]

Outline

- Overview of Higgs bundles on local CY_3 , G_2 and $Spin(7)$ spaces

M-theory/Type IIA perspective $\rightarrow D=5,4,3$ $N=1$

[Also F-theory/IIB on elliptically fibered $CY_4 \rightarrow D=4$ $N=1$]

- $Spin(7)$ Higgs bundles and relation to G_2 and CY_4 Higgs bundles
 \rightarrow geometric unification
- Zero modes of $Spin(7)$ Higgs bundles – classical results

Quantum results c.f., J. Heckman's talk

- Conclusions/Outlook

Overview of Higgs bundles

Within M-theory (11D) compactification: localized gauge degrees live on co-dimension four ADE singular subspace of a special holonomy spaces:

ALE fibration

$$\begin{array}{l} \mathbb{R}^4/\Gamma_{ADE} \rightarrow X_6 \\ \downarrow \\ \Sigma_2 \end{array} \quad \begin{array}{l} \text{Local Calabi-Yau threefold} \\ \\ \text{2D Riemann surface } \Sigma_2 \end{array} \quad \rightarrow \text{D=5 N=1}$$

$$\begin{array}{l} \mathbb{R}^4/\Gamma_{ADE} \rightarrow X_7 \\ \downarrow \\ Q_3 \end{array} \quad \begin{array}{l} \text{Local } G_2 \text{ holonomy space} \\ \\ \text{3D subspace } Q_3 \end{array} \quad \rightarrow \text{D=4 N=1}$$

$$\begin{array}{l} \mathbb{R}^4/\Gamma_{ADE} \rightarrow X_8 \\ \downarrow \\ M_4 \end{array} \quad \begin{array}{l} \text{Local Spin(7) holonomy space} \\ \\ \text{4D subspace } M_4 \end{array} \quad \rightarrow \text{D=3 N=1}$$

ALE subspaces as gauge sectors localized on D6-brane

Singular codimension four subspaces can be interpreted in Type IIA string theory as D6 branes wrapping respective singular subspaces.

[via duality: M-theory on $S^1 \sim$ Type IIA string theory]

- Start with 10D Super Yang-Mills (SYM) action
w/adjoint valued gauge fields A_μ & fermions ψ^α (Majorana-Weyl)

$$S = \int d^{10}x \text{Tr} \left(-\frac{1}{2} F^{\mu\nu} F_{\mu\nu} + i\psi^\alpha \gamma_{\alpha\beta}^\mu D_\mu \psi^\beta \right)$$

$$D_\mu = \partial_\mu + gA_\mu$$
$$F_{\mu\nu} = [D_\mu, D_\nu]$$

& supersymmetry transformations

ALE subspaces as gauge sectors localized on D6-brane

- **Truncation to 7D:** D6-brane worldvolume action
w/ gauge fields A_μ , 3 adjoint rep. scalars ϕ_i (3 "transverse dims")
& pair of fermions ψ_i (w/symplectic Majorana constraint)

$$S = \int dt d^6x \operatorname{tr} \left[-\frac{1}{4} F_{\mu\nu} F^{\mu\nu} - \frac{1}{4} D_\mu \phi_a D^\mu \phi_a + \frac{1}{4} [\phi_a, \phi_b] [\phi_a, \phi_b] - \frac{i}{2} \bar{\Psi}^I \not{D} \Psi_I - \frac{i}{2} \sigma_{IJ}^a \bar{\Psi}^I [\phi_a, \Psi^J] \right]$$

& supersymmetry transformations

R-symmetry
(3 transverse dims)

Global symmetry: $\operatorname{Spin}(1,9) \rightarrow \operatorname{Spin}(1,6) \otimes \operatorname{Spin}(3)_R$

- **Compactify** D6-brane worldvolume action on Σ_2, Q_3, M_4 spaces \rightarrow
Higgs bundles

D6-brane on Σ_2 , Q_3 , M_4 spaces \rightarrow Higgs bundles

Global symmetry & Topological twist

on Σ_2

$$Spin(1,6) \otimes Spin(3)_R \rightarrow Spin(1,4) \otimes Spin(2) \otimes Spin(3)_R$$

$$\simeq U(1) \quad \swarrow \quad \searrow$$

$$U(1)_{\text{twist}}$$

Diagonal $U(1) \rightarrow$ twist
preserves D=5 N=1 SUSY

on Q_3

$$Spin(1,6) \otimes Spin(3)_R \rightarrow Spin(1,3) \otimes Spin(3) \otimes Spin(3)_R$$

$$\searrow \quad \swarrow$$

$$SU(2)_{\text{twist}}$$

Diagonal $SU(2) \rightarrow$ twist
preserves D=4 N=1 SUSY

on M_4

$$Spin(1,6) \otimes Spin(3)_R \rightarrow Spin(1,2) \otimes Spin(4) \otimes Spin(3)_R$$

$$\simeq SU(2)_l \otimes SU(2)_r \quad \swarrow \quad \searrow$$

$$SU(2)_{\text{twist}}$$

Diagonal $SU(2) \rightarrow$ twist
preserves D=3 N=1 SUSY

Compactification on Σ_2

- **Higgs bundle:** 2D SUSY gauge theory on Riemann surface Σ_2 w/ one-form Higgs field ϕ_{Hit} & vector bundle V connection A (in adjoint reps. of ADE groups)

N. Hitchin, Lon. Math. Soc. 55 (1987)

- **Hitchin equations:** $\bar{\partial}_A \phi_{\text{Hit}} = 0$
(BPS eqs.)
$$F_A + \frac{i}{2} [\phi_{\text{Hit}}^\dagger, \phi_{\text{Hit}}] = 0$$

Unitary frame [\dagger - hermitian conjugate]

Solutions specified modulo gauge transformations

- ϕ_{Hit} characterize local deformations in the shape of the ambient Calabi-Yau threefold
 A parameterizes gauge field deformations (flux), sourced by D-brane

Hitchin System

Compactification on Q₃

[Pantev, Wijnholt '09] [Acharya] [Barbosa '19]

- **Higgs bundle:** 3D gauge SUSY gauge theory on Q₃
adjoint valued one-form Higgs field ϕ , vector bundle connection A
(after twisting)
- **BPS equations:** $D_A \phi = 0$ $D_A \star \phi = 0$
 $F = [\phi, \phi]$

Complexified connection: $\mathcal{A} = A + i\phi \rightarrow$

BPS equations: $\mathcal{F} = 0$ $D_A \star \phi = 0$
[$D_A \phi = 0$ redundant]

Pantev Wijnholt (PW) System

Compactification on Q_3

[Pantev, Wijnholt '09] [Acharya] [Barbosa '19]

- **Chiral matter** localized at the center of solitonic configuration of ϕ , with vanishing value there (co-dimension seven singularities).

[Braun, Cizel, Hübner, Schäfer-Nameki '18]
[Hübner'20]

- Subsequent studies developed and extended to **co-dimension six singularities** (non-chiral matter). Appealing feature: they could possibly connect to building **compact G_2** manifolds via **twisted connected sums (TCS)**.

[Kovalev '03] [Corti, Haskins, Nordström, Pacini '13-'15]
[Braun, Schäfer-Nameki '18]

- Most prior analyses of localized matter have assumed one-form ϕ is diagonal and no A in Q_3 

[Barbosa, M.C., Heckman, Lawrie, Torres, Zoccarato '19]
c.f., M.C.'s talk at the Santa Barbara Collaboration Meeting '19

PW system generalized non-Abelian case & to include non-diagonal one-form ϕ & non-zero flux A

$$\phi \sim \begin{bmatrix} \lambda & 1 & 0 \\ z & \lambda & 0 \\ 0 & 0 & -2\lambda \end{bmatrix}$$

- One-form ϕ components will not commute
- ``T-brane type configuration''
(fits in the broader scheme of T-brane like phenomena)
- Local model: three-manifold Q_3 as a Riemann surface Σ_2 fibered over an interval S^1 : the gauge theory on Σ_2 is a *mild deformation* of a Hitchin system on Σ_2
- As Hitchin system on Σ_2 describes a local Calabi-Yau geometry \rightarrow obtain a local deformation of a TCS-like construction \rightarrow Interpreted as building up a local G_2 background
- Localized modes (in co-dim seven)
 \rightarrow explicit solutions; their existence reduces to a linear algebraic problem
[à la localized zero modes of T-branes F-theory on Calabi-Yau fourfolds]
[Cecotti, Cheng, Heckman, Vafa'09]
[Cecotti, Cordova, Heckman, Vafa'10]

[Spin(7); D=3 N=1]

Compactification on M_4

[Vafa, Witten '94]

[Heckman, Lawrie, Ling, Zoccarato '18]

[M.C., Heckman, Rochais, Torres, Zoccarato '20]

- **Higgs bundle: 4D gauge theory on M_4**
Higgs field ϕ - an adjoint valued **self-dual two-form** &
gauge field strength F_{SD} - **self-dual** part of the bundle curvature

- **BPS equations:** $D_A \phi = 0$

$$F_{SD} + \phi \times \phi = 0$$

Cross product: $(\phi \times \phi)_{ij} = \frac{1}{4}[\phi_{ik}, \phi_{jl}]g^{kl}$

Equations those of N=4 SYM on non-Kähler M_4 w/ topological twist

[Vafa, Witten '94]

Vafa-Witten (VW) System

VW system includes other Higgs bundles as solutions



Geometric unification of Higgs bundle vacua

- PW system related to VW system:

Consider VW system on $M_4 = Q_3 \times S^1$ S^1 parameterized by t

Write the SD forms as: $\phi = \hat{\phi} \wedge dt + \star_3 \hat{\phi}$ $\hat{\phi} \in \Omega^1(Q_3)$

- BPS equations:** $F - [\hat{\phi}, \hat{\phi}] + \star_3 (D_t A - d_3 A_t) = 0$ $D_A \hat{\phi} + \star_3 D_t \hat{\phi} = 0$

$$D_A \star_3 \hat{\phi} = 0$$

- One recovers PW w/ $A_t = \partial_t A = \partial_t \hat{\phi} = 0$

PW is the dimensional reduction of VW along S^1

Digression: another way to obtain D=4 N=1

Within F-theory (12D) compactification: localized gauge degrees live on co-dimension four ADE singular subspace of elliptically fibered Calabi-Yau fourfold

$$\begin{array}{ccc} \mathbb{R}^4/\Gamma_{ADE} \rightarrow X_8 & \text{Local elliptically fibered Calabi-Yau fourfold} & \\ \downarrow & & \\ S_4 & \text{Kähler 4D subspace} & \text{D=4 N=1} \end{array}$$

- Singular co-dim 4 subspace S_4 can be interpreted as D7 brane wrapping S_4 [duality: F-theory on $T^2 \rightarrow$ Type IIB string theory]

\rightarrow 8D supersymmetric gauge theory of D7-brane worldvolume action, compactified on S_4



[Beasley, Heckman, Vafa '08]

Beasley Heckman Vafa (BHV) system

[Elliptically fibered CY fourfold; D=4 N=1]

BHV system

[Beasley, Heckman, Vafa '08]

- **Higgs bundle** on a Kähler 4D manifold S_4
Higgs field is a complex adjoint valued (2,0)-form (after twisting)

- BPS equations:

$$\bar{\partial}_A \phi = 0$$

$$F_{(0,2)} = 0$$

Kähler form


$$J \wedge F + \frac{i}{2} [\phi^\dagger, \phi] = 0$$

Holomorphic structure makes construction of solutions feasible
(even non-Abelian ones)

Geometric unification of Higgs bundle vacua

- **BHV system is also related to VW system**

- Take VW system with $M_4 \sim S_4$ to be Kähler manifold
- SD two-forms admit a Hodge decomposition

$$\Omega_+^2(S_4) \simeq H^{(2,0)}(S_4) \oplus H^{(0,2)}(S_4) \oplus H_{n.p.}^{(1,1)}(S_4)$$

← non-primitive

- VW equations become:

$$F_{(0,2)} - \frac{i}{2} \phi_{(1,1)} \times \phi_{(0,2)}^\dagger = 0 \qquad \bar{\partial}_A \phi_{(2,0)} - \frac{i}{2} \partial_A \phi_{(1,1)} = 0$$
$$J \wedge F = \frac{i}{2} [\phi_{(2,0)}, \phi_{(0,2)}^\dagger]$$

BHV is recovered by setting (1,1)-component set to zero

Geometric unification of Higgs bundle vacua

$$BHV \subset VW \supset PW$$

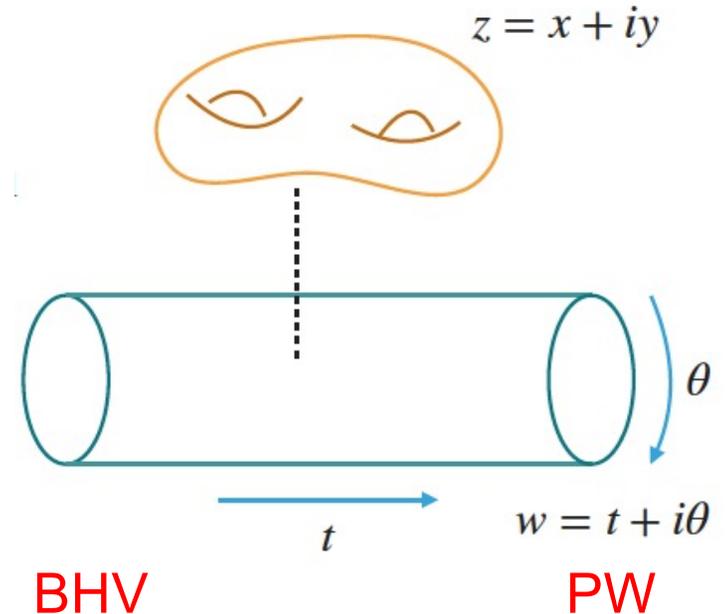
- PW to BHV interpolation:**

Take M_4 manifold: $M_4 = \Sigma_2 \times \mathbb{R} \times S^1$
 $(x,y) \quad t \quad \theta$

$$|\phi_{(1,1)}| \sim e^{\lambda t} \quad t \rightarrow -\infty$$

$$|\partial_\theta \psi| \sim e^{-\lambda_2 t} \quad |A_\theta| \sim e^{-\lambda_1 t} \quad t \rightarrow +\infty$$

- For $t \rightarrow -\infty$ solution approaches BHV
- For $t \rightarrow +\infty$ solution approaches PW
- In the middle there is a Spin(7) solution
(Explicit construction of the full solution)



Higgs bundle version of TCS construction of Spin(7) manifolds

[Braun, Schäfer-Nameki '18]

- Also examples w/ PW to PW interpolation:** $M_4 = Q_3 \times \mathbb{R} \quad Q_3 = \mathbf{T}^3$

Zero Modes of VW system – classical

- Zero modes equations – infinitesimal variation around a solution:

$$A = \langle A \rangle + a \quad \phi = \langle \phi \rangle + \varphi$$

- Zero mode equations for VW:

[For simplicity: abelian, without flux $\langle A \rangle = 0$]

$$D_A a + \star D_A a + \phi \times \varphi = 0$$

$$D_A \varphi - [\phi, a] = 0$$

drop $\langle \rangle$ for $\langle \phi \rangle$

Solutions identified via inf. gauge transformations: $a \simeq a + D_A \xi$
 $\varphi \sim \varphi + [\phi, \xi]$

- Where are zero modes localized?
[loci where $\phi=0$ (due to $\phi \rightarrow t \phi$ rescaling)]

Can we localize zero modes on M_4 ?

As $BHV \subset VW \supset PW$, one expects that localized VW zero modes will share some similarities with PW and/or BHV ones.

- **Comparison to PW:** three scalars, generically localize modes at co-dim three subspace of Q_3 , i.e. a point.

VW - three independent components of SD scalars, to localize modes in co-dim three in M_4 , i.e. **a line**

- **Comparison to BHV:** on a Kähler surface S_4 , w/ one (2,0)-form, w/ zeroes localized in complex co-dim one subspace. Gauge **flux** further localizes to complex co-dim two in S_4 , i.e. a point.

VW – consider a special case of BHV background, to localize in real codimension four on M_4 , i.e. **a point**

- **Full-fledged VW:** zero modes generically lie in codimension three in M_4 , i.e. **a line**

[Pull-back of flux on line does not produce further localization]

Supported by **explicit analysis in a local patch** $M_4 = \mathbb{R}^4$

[Clash of PW and BHV complex structures]

Quantum corrections could lift zero modes

[M.C., Heckman, Torres, Zoccarato '21]

In spite of lack of the power of holomorphicity (unlike $D=4$ $N=1$), there is certain quantum protection in $D=3$ $N=1$, due to

- Discrete Symmetry Anomalies involving space-time reflection symmetries, inherited from 7D SYM action (D6-brane worldvolume) reduced on M_4
- Examples of gluing local 4D patches into compact M_4 resulting in $D=3$ $N=1$ & counting of quantum protected zero modes

c.f., J. Heckman's talk

Conclusions

- Higgs bundles in Spin(7) compactifications of M-theory
- This system provides a unification of other known Higgs bundles (contains PW and BHV bundles as special cases)
- Build solutions interpolating between BHV and PW & interfaces between PW solutions
- Localization of zero modes: generically localized at most on codimension three in M_4

Classical results \rightarrow quantum ones

c.f., J. Heckman's talk

Outlook

- Most of the analysis on (ultra-)local patches
→ full-fledged global analysis:
 - Higgs bundles on compact M_4
Embedding in non-compact $\text{Spin}(7)$
 - Embedding in compact $\text{Spin}(7)$ [D=3 N=1 w/ gravity]

Progress c.f., J. Heckman's talk

- Toward geometric classification of D=3 N=1 dualities
- Type IIA on $\text{Spin}(7)$ D=2 N=(1,1)
Type IIB on $\text{Spin}(7)$ D=2 N=(0,2) – holomorphicity?

Thank you!