

# Solitons in Bryant's $G_2$ -Laplacian flow.

Mark Haskins  
Duke University

9th September 2021  
2021 Simons Collaboration on Special Holonomy  
in Geometry, Analysis and Physics Annual Meeting

Joint work in progress with  
Rowan Juneman & Johannes Nordström (Bath)

## 3-forms in 7 variables and $G_2$ -structures

---

One way to define  $G_2$  is as  $G_2 = \text{Aut}(\mathbb{O})$  where  $\mathbb{O}$  is the octonions.

Define a cross-product and a  $G_2$ -**invariant** 3-form  $\varphi_0$  on  $\mathbb{R}^7 = \text{Im}(\mathbb{O})$  using octonionic multiplication and the Euclidean inner product

$$u \times v := \text{Im}(uv)$$

$$\varphi_0(u, v, w) := \langle u \times v, w \rangle = \langle uv, w \rangle.$$

## 3-forms in 7 variables and $G_2$ -structures

---

One way to define  $G_2$  is as  $G_2 = \text{Aut}(\mathbb{O})$  where  $\mathbb{O}$  is the octonions.

Define a cross-product and a  $G_2$ -**invariant** 3-form  $\varphi_0$  on  $\mathbb{R}^7 = \text{Im}(\mathbb{O})$  using octonionic multiplication and the Euclidean inner product

$$u \times v := \text{Im}(uv)$$
$$\varphi_0(u, v, w) := \langle u \times v, w \rangle = \langle uv, w \rangle.$$

For an oriented smooth 7-manifold  $M$  and  $p \in M$

$$\mathcal{P}_p(M) := \{\varphi \in \Lambda^3 T_p^* M \mid \iota^* \varphi_0 = \varphi \text{ for } \iota: T_p M \rightarrow \mathbb{R}^7\}$$

where  $\iota$  is any orientation-preserving isomorphism.

A 3-form  $\varphi$  on  $M$  is *positive* if  $\varphi$  is a section of  $\mathcal{P}(M)$ , i.e.  $\varphi_p \in \mathcal{P}_p(M) \forall p$ , where  $\mathcal{P}(M)$  denotes the bundle over  $M$  with fibre  $\mathcal{P}_p(M)$ .

## 3-forms in 7 variables and $G_2$ -structures

---

One way to define  $G_2$  is as  $G_2 = \text{Aut}(\mathbb{O})$  where  $\mathbb{O}$  is the octonions.

Define a cross-product and a  $G_2$ -invariant 3-form  $\varphi_0$  on  $\mathbb{R}^7 = \text{Im}(\mathbb{O})$  using octonionic multiplication and the Euclidean inner product

$$u \times v := \text{Im}(uv)$$
$$\varphi_0(u, v, w) := \langle u \times v, w \rangle = \langle uv, w \rangle.$$

For an oriented smooth 7-manifold  $M$  and  $p \in M$

$$\mathcal{P}_p(M) := \{\varphi \in \Lambda^3 T_p^* M \mid \iota^* \varphi_0 = \varphi \text{ for } \iota: T_p M \rightarrow \mathbb{R}^7\}$$

where  $\iota$  is any orientation-preserving isomorphism.

A 3-form  $\varphi$  on  $M$  is *positive* if  $\varphi$  is a section of  $\mathcal{P}(M)$ , i.e.  $\varphi_p \in \mathcal{P}_p(M) \forall p$ , where  $\mathcal{P}(M)$  denotes the bundle over  $M$  with fibre  $\mathcal{P}_p(M)$ .

Each positive 3-form on  $M$  defines a reduction of the frame bundle  $\mathcal{F}M$  to a principal subbundle of  $\mathcal{F}M$  with fibre  $G_2$ , i.e. a  $G_2$ -structure on  $M$  that induces the given orientation on  $M$ .

Positive 3-forms on  $M \iff$  (oriented)  $G_2$ -structures on  $M$ .

# 1st-order PDE system for $G_2$ holonomy metrics

---

Holonomy/parallel tensors correspondence:  $\text{Hol}_g(M) \subseteq G_2 \subset \text{SO}(7)$  implies

$M^7$  admits a  $g$ -parallel positive 3-form  $\varphi$ .

Converse?

**Theorem:** Let  $(M, \varphi, g_\varphi)$  be a  $G_2$ -structure; the following are equivalent

1.  $\text{Hol}(g_\varphi) \subseteq G_2$  and  $\varphi$  is the induced 3-form
  2.  $d\varphi = d^*\varphi = 0$ , where  $d^*$  is defined using Hodge star  $*$  w.r.t.  $g_\varphi$ .
- 2 is *nonlinear* in  $\varphi$ :  $g_\varphi$  depends nonlinearly on  $\varphi$  and  $d^*$  depends on  $g_\varphi$ .

# 1st-order PDE system for $G_2$ holonomy metrics

---

**Theorem:** Let  $(M, \varphi, g_\varphi)$  be a  $G_2$ -structure; the following are equivalent

1.  $\text{Hol}(g_\varphi) \subseteq G_2$  and  $\varphi$  is the induced 3-form
2.  $d\varphi = d^*\varphi = 0$ , where  $d^*$  is defined using Hodge star  $*$  w.r.t.  $g_\varphi$ .

2 is *nonlinear* in  $\varphi$ :  $g_\varphi$  depends nonlinearly on  $\varphi$  and  $d^*$  depends on  $g_\varphi$ .

By writing equation for 3-form  $\varphi$  (not metric  $g$  directly) and allowing  $\text{Hol}(g_\varphi) \subseteq G_2$  we obtain a *PDE system*: (not integro-differential equations) a **1st-order system of  $49=(35+21-7)$  equations on the 35 coeffs of  $\varphi$ !** It is an overdetermined diffeomorphism-invariant system.

# 1st-order PDE system for $G_2$ holonomy metrics

---

**Theorem:** Let  $(M, \varphi, g_\varphi)$  be a  $G_2$ -structure; the following are equivalent

1.  $\text{Hol}(g_\varphi) \subseteq G_2$  and  $\varphi$  is the induced 3-form
2.  $d\varphi = d^*\varphi = 0$ , where  $d^*$  is defined using Hodge star  $*$  w.r.t.  $g_\varphi$ .

**2** is *nonlinear* in  $\varphi$ :  $g_\varphi$  depends nonlinearly on  $\varphi$  and  $d^*$  depends on  $g_\varphi$ .

By writing equation for 3-form  $\varphi$  and allowing  $\text{Hol}(g_\varphi) \subseteq G_2$  we obtain a *PDE system*: a **1st-order system of 49=(35+21-7) equations on the 35 coeffs of  $\varphi$ !** It is an overdetermined diffeomorphism-invariant system.

**Elliptic approach:** (singular) perturbation method due to **Joyce**.

**All** known constructions of  $G_2$ -holonomy metrics on compact manifolds use this approach!

*Idea:* Construct a *closed*  $G_2$ -structure  $\varphi$  with  $d^*\varphi$  *sufficiently small*. Condition that closed 3-form  $\tilde{\varphi} = \varphi + d\eta$  solves **2** yields a **nonlinear elliptic PDE for 2-form  $\eta$**  which is solvable by an iteration method ( $\eta$  is small).

Difficulty then becomes to construct initial closed  $G_2$ -structure  $\varphi$  with sufficiently small torsion: all constructions exploit various *degenerate limits*.

## Bryant's Laplacian flow

---

There is a natural flow on *closed*  $G_2$ -structures. Solve

$$\frac{d\varphi_t}{dt} = \Delta_{\varphi_t}\varphi_t \quad (\text{LF})$$

with initial condition  $\varphi_0$  satisfying  $d\varphi_0 = 0$ . (Then  $d\varphi_t = 0$  for all  $t$ .)



## Bryant's Laplacian flow

---

There is a natural flow on *closed*  $G_2$ -structures. Solve

$$\frac{d\varphi_t}{dt} = \Delta_{\varphi_t}\varphi_t \quad (\text{LF})$$

with initial condition  $\varphi_0$  satisfying  $d\varphi_0 = 0$ .

- Induced metric  $g_t$  evolves under (LF) by

$$\frac{dg_t}{dt} = -2\text{Ric}(g_t) + \text{terms quadratic in torsion of } \varphi_t$$

## Bryant's Laplacian flow

---

There is a natural flow on *closed*  $G_2$ -structures. Solve

$$\frac{d\varphi_t}{dt} = \Delta_{\varphi_t} \varphi_t \quad (\text{LF})$$

with initial condition  $\varphi_0$  satisfying  $d\varphi_0 = 0$ .

- Induced metric  $g_t$  evolves under (LF) by

$$\frac{dg_t}{dt} = -2\text{Ric}(g_t) + \text{terms quadratic in torsion of } \varphi_t$$

- Stationary points of (LF) are exactly **torsion-free  $G_2$ -structures**.
- (LF) is the (upward) gradient flow for **Hitchin's volume functional**  $\text{vol}(\varphi)$

$$\text{vol}(\varphi) := \frac{1}{7} \int_M \varphi \wedge * \varphi$$

when restricted to cohomology class of  $\varphi_0$ . Critical points of  $\text{vol}(\varphi)$  in  $[\varphi]$  are **maxima** (strict modulo diffeos).

## Bryant's Laplacian flow

---

There is a natural flow on *closed*  $G_2$ -structures. Solve

$$\frac{d\varphi_t}{dt} = \Delta_{\varphi_t}\varphi_t \quad (\text{LF})$$

with initial condition  $\varphi_0$  satisfying  $d\varphi_0 = 0$ .

- Induced metric  $g_t$  evolves under (LF) by

$$\frac{dg_t}{dt} = -2\text{Ric}(g_t) + \text{terms quadratic in torsion of } \varphi_t$$

- Stationary points of (LF) are exactly **torsion-free  $G_2$ -structures**.
- (LF) is the (upward) gradient flow for **Hitchin's volume functional**  $\text{vol}(\varphi)$  when restricted to cohomology class of  $\varphi_0$ . Critical points of  $\text{vol}(\varphi)$  in  $[\varphi]$  are **maxima** (strict modulo diffeos).
- On a compact manifold  $\text{vol}(\varphi_t)$  is *increasing* along (LF)  
 $\Rightarrow$  there are no *compact* shrinking solitons in (LF) (unlike in Ricci flow).

## Bryant's Laplacian flow

---

There is a natural flow on *closed*  $G_2$ -structures. Solve

$$\frac{d\varphi_t}{dt} = \Delta_{\varphi_t}\varphi_t \quad (\text{LF})$$

with initial condition  $\varphi_0$  satisfying  $d\varphi_0 = 0$ .

- Induced metric  $g_t$  evolves under (LF) by

$$\frac{dg_t}{dt} = -2\text{Ric}(g_t) + \text{terms quadratic in torsion of } \varphi_t$$

- Stationary points of (LF) are exactly **torsion-free  $G_2$ -structures**.
- (LF) is the (upward) gradient flow for **Hitchin's volume functional**  $\text{vol}(\varphi)$  when restricted to cohomology class of  $\varphi_0$ . Critical points of  $\text{vol}(\varphi)$  in  $[\varphi]$  are **maxima** (strict modulo diffeos).
- On a compact manifold  $\text{vol}(\varphi_t)$  is *increasing* along (LF)
- Bryant–Xu : Short-time existence & uniqueness of solutions to (LF)
- Lotay–Wei: Torsion-free  $G_2$ -structures are stable under (LF).  
Lotay–Wei also establish analogues of results in Ricci flow, e.g. Hamilton's compactness result, Shi estimates, finite-time extension result.

## Geometric flows 101

---

- *Ideal case*: Establish **long-time existence** & **convergence of flow** as  $t \rightarrow \infty$  to a static solution.

e.g. *Hamilton's 3-diml spherical space form theorem*: any initial metric  $g_0$  on a compact 3-manifold with  $Ric(g_0) > 0$  converges under normalized Ricci flow to an Einstein metric  $Ric(g) = g$  (which must be a spherical space form in 3 dims).

Later influential higher diml results under other *curvature pinching conditions* on initial metric: Huisken, Hamilton, Böhm-Wilking, Brendle-Schoen.

## Geometric flows 101

---

- *Ideal case*: Establish **long-time existence** & **convergence of flow** as  $t \rightarrow \infty$  to a static solution.

e.g. *Hamilton's 3-diml spherical space form theorem*: any initial metric  $g_0$  on a compact 3-manifold with  $Ric(g_0) > 0$  converges under normalized Ricci flow to an Einstein metric  $Ric(g) = g$  (which must be a spherical space form in 3 dims).

- In general **long-time existence** or **convergence to a static solution** may fail. Both issues arise in 3-dim Ricci flow for *general* initial metrics.

**Long-time existence** fails when we encounter **finite-time singularities** of the flow, e.g. in Ricci flow when  $|Riem(g_t)| \rightarrow \infty$  as  $t \nearrow T$ .

# Geometric flows 101

---

- *Ideal case*: Establish **long-time existence** & **convergence of flow** as  $t \rightarrow \infty$  to a static solution.

e.g. *Hamilton's 3-diml spherical space form theorem*: any initial metric  $g_0$  on a compact 3-manifold with  $Ric(g_0) > 0$  converges under normalized Ricci flow to an Einstein metric  $Ric(g) = g$  (which must be a spherical space form in 3 dims).

- In general **long-time existence** or **convergence to a static solution** may fail.

**Long-time existence** fails when we encounter **finite-time singularities** of the flow, e.g. in Ricci flow when  $|Riem(g_t)| \rightarrow \infty$  as  $t \nearrow T$ .

In 3-diml Ricci flow *some* singularities were expected/needed:

*Thurston's Geometrization Conjecture* provided important insights into what (finite and infinite time) singularities *should* arise in 3-dim Ricci flow.

e.g. **Neck-pinch singularities** implementing *connect sum decomposition* of initial 3-manifold; modelled on **shrinking cylinder**  $S^2 \times \mathbb{R}$  which is a simple example of a **shrinking soliton**.

# Geometric flows 101

---

- *Ideal case*: Establish **long-time existence** & **convergence of flow** as  $t \rightarrow \infty$  to a static solution.

e.g. *Hamilton's 3-diml spherical space form theorem*: any initial metric  $g_0$  on a compact 3-manifold with  $Ric(g_0) > 0$  converges under normalized Ricci flow to an Einstein metric  $Ric(g) = g$  (which must be a spherical space form in 3 dims).

- In general **long-time existence** or **convergence to a static solution** may fail.

**Long-time existence** fails when we encounter **finite-time singularities** of the flow,

e.g. in Ricci flow when  $|Riem(g_t)| \rightarrow \infty$  as  $t \nearrow T$ .

e.g. **Neck-pinch singularities** implementing *connect sum decomposition* of initial 3-manifold; modelled on **shrinking cylinder**  $S^2 \times \mathbb{R}$  which is a simple example of a **shrinking soliton**.

In general need to develop a theory of **finite-time singularity models**:

*if* only 'expected' finite-time singularities arise can hope to develop a **flow with surgeries** and study its long-time behaviour, eg *Ricci flow with surgery*.

We don't have any analogue of the Geometrization Conjecture to guide us in the  $G_2$  setting but we can study **solitons in the Laplacian flow**.



# Solitons in the Laplacian flow

---

In many geometric flows **solitons** provide key models for **singularity formation** (and sometimes **singularity resolution**).

A **Laplacian soliton** is a  $G_2$ -structure  $\varphi$ , vector field  $X$ ,  $\lambda \in \mathbb{R}$  satisfying

$$\begin{cases} d\varphi = 0 \\ \Delta_\varphi \varphi = \lambda\varphi + \mathcal{L}_X \varphi \end{cases} \quad (\text{LSE})$$

$\Leftrightarrow$  self-similar solution of Laplacian flow

$$\varphi_t = k(t)^3 f^* \varphi, \quad \frac{df}{dt} = k(t)^{-2} X, \quad k(t) = \frac{1}{3}(3 + 2\lambda t)$$

$\lambda < 0$ : **shrinkers** – **ancient solutions**, i.e. exist backwards to  $t = -\infty$

$\lambda = 0$ : **steady solitons** – **eternal solutions**, i.e. exist for all time  $t \in \mathbb{R}$

$\lambda > 0$ : **expanders** – **immortal solutions**, i.e. exist up to  $t = +\infty$

# Solitons in the Laplacian flow

---

In many geometric flows **solitons** provide key models for **singularity formation** (and sometimes **singularity resolution**).

A **Laplacian soliton** is a  $G_2$ -structure  $\varphi$ , vector field  $X$ ,  $\lambda \in \mathbb{R}$  satisfying

$$\begin{cases} d\varphi = 0 \\ \Delta_\varphi \varphi = \lambda\varphi + \mathcal{L}_X \varphi \end{cases} \quad (\text{LSE})$$

$\Leftrightarrow$  self-similar solution of Laplacian flow

$$\varphi_t = k(t)^3 f^* \varphi, \quad \frac{df}{dt} = k(t)^{-2} X, \quad k(t) = \frac{1}{3}(3 + 2\lambda t)$$

$\lambda < 0$ : **shrinkers** – **ancient solutions**, i.e. exist backwards to  $t = -\infty$

$\lambda = 0$ : **steady solitons** – **eternal solutions**, i.e. exist for all time  $t \in \mathbb{R}$

$\lambda > 0$ : **expanders** – **immortal solutions**, i.e. exist up to  $t = +\infty$

- Non-steady soliton  $\Rightarrow \varphi$  exact
- Solitons on a *compact* manifold are stationary or expanders
- Scaling behaviour:  $(\varphi, X)$  a  $\lambda$ -soliton  $\iff (k^3\varphi, k^{-2}X)$  a  $k^{-2}\lambda$ -soliton.
- Bryant (unpublished) has studied the local generality of Laplacian solitons using methods from overdetermined systems of PDE.

## Laplacian soliton construction methods

---

No general analytic machinery to construct *global* solutions to (LSE). Have same problem in Ricci flow (except in the Kähler setting.) Must study Laplacian solitons with *additional geometric structure*.

- *Homogeneous solitons*: Quite a bit of work by Lauret, Fino & coworkers  
Work has a Lie-theoretic flavour – nilpotent and solvable Lie groups  
Relevance to finite-time singularity formation on say compact 1-connected manifolds  $M$  with  $p_1(M) \neq 0$  is unclear.

## Laplacian soliton construction methods

---

No general analytic machinery to construct *global* solutions to (LSE). Must study Laplacian solitons with *additional geometric structure*.

- *Homogeneous solitons*: Quite a bit of work by Lauret, Fino & coworkers  
Work has a Lie-theoretic flavour – nilpotent and solvable Lie groups  
Relevance to finite-time singularity formation on say compact 1-connected manifolds  $M$  with  $p_1(M) \neq 0$  is unclear.
- *Cohomogeneity-one solitons*: a Lie group acts with codimension-one generic orbit; reduces (LSE) to nonlinear system of ODEs.  
Various important Ricci solitons are cohomogeneity one: Cigar soliton; shrinking round cylinders; Bryant's steady soliton and expanders; Cao's Kähler expander and steady soliton; FIK Kähler shrinkers

## Laplacian soliton construction methods

---

No general analytic machinery to construct *global* solutions to (LSE). Must study Laplacian solitons with *additional geometric structure*.

- *Homogeneous solitons*: Quite a bit of work by Lauret, Fino & coworkers  
Work has a Lie-theoretic flavour – nilpotent and solvable Lie groups  
Relevance to finite-time singularity formation on say compact 1-connected manifolds  $M$  with  $p_1(M) \neq 0$  is unclear.
- **Cohomogeneity-one solitons**: a Lie group acts with codimension-one generic orbit; reduces (LSE) to nonlinear system of ODEs.
- *Bundle constructions*, e.g. a circle/torus bundle over base manifold with special geometry. Often leads to spaces foliated by special hypersurfaces and governed by the same ODEs as some cohomogeneity one examples.

## Laplacian soliton construction methods

---

No general analytic machinery to construct *global* solutions to (LSE). Must study Laplacian solitons with *additional geometric structure*.

- *Homogeneous solitons*: Quite a bit of work by Lauret, Fino & coworkers  
Work has a Lie-theoretic flavour – nilpotent and solvable Lie groups  
Relevance to finite-time singularity formation on say compact 1-connected manifolds  $M$  with  $p_1(M) \neq 0$  is unclear.
- *Cohomogeneity-one solitons*: a Lie group acts with codimension-one generic orbit; reduces (LSE) to nonlinear system of ODEs.
- *Bundle constructions*, e.g. a circle/torus bundle over base manifold with special geometry.
- *Solitons with special torsion*: specific to  $G_2$  case. Gavin Ball studied closed  $G_2$ -structures with special torsion in his Duke thesis and found some steady Laplacian solitons this way.

# Laplacian soliton construction methods

---

No general analytic machinery to construct *global* solutions to (LSE). Must study Laplacian solitons with *additional geometric structure*.

- *Homogeneous solitons*: Quite a bit of work by Lauret, Fino & coworkers  
Work has a Lie-theoretic flavour – nilpotent and solvable Lie groups  
Relevance to finite-time singularity formation on say compact 1-connected manifolds  $M$  with  $p_1(M) \neq 0$  is unclear.
- *Cohomogeneity-one solitons*: a Lie group acts with codimension-one generic orbit; reduces (LSE) to nonlinear system of ODEs.
- *Bundle constructions*, e.g. a circle/torus bundle over base manifold with special geometry.
- *Solitons with special torsion*: specific to  $G_2$  case. Gavin Ball found some steady Laplacian solitons this way.
- *$S^1$ -collapsed solitons?*: Foscolo-H-Nordström (accepted **Duke Math J**) constructed many new complete noncompact torsion-free  $G_2$ -structures on circle bundles over AC Calabi-Yau 3-folds. ?  $\exists$  construction of  $S^1$ -collapsed solitons that is related to solitons of Lagrangian MCF in CY 3-folds?

# Laplacian soliton construction methods

---

No general analytic machinery to construct *global* solutions to (LSE). Must study Laplacian solitons with *additional geometric structure*.

- *Homogeneous solitons*: Quite a bit of work by Lauret, Fino & coworkers  
Work has a Lie-theoretic flavour – nilpotent and solvable Lie groups  
Relevance to finite-time singularity formation on say compact 1-connected manifolds  $M$  with  $p_1(M) \neq 0$  is unclear.
- *Cohomogeneity-one solitons*: a Lie group acts with codimension-one generic orbit; reduces (LSE) to nonlinear system of ODEs.
- *Bundle constructions*, e.g. a circle/torus bundle over base manifold with special geometry.
- *Solitons with special torsion*: specific to  $G_2$  case. Gavin Ball found some steady Laplacian solitons this way.
- *$S^1$ -collapsed solitons?*: ?  $\exists$  construction of  $S^1$ -collapsed solitons that is related to solitons of Lagrangian MCF in CY 3-folds?  
Minimal symmetry assumption, but not clear that highly-collapsed solitons could appear as finite-time singularity models. But (as in torsion-free case) might also suggest existence of new non-collapsed solitons.



# Our motivation and high-level overview of results

---

**Motivation:** Construct cohomogeneity one Laplacian solitons that are potentially viable models for finite-time singularity formation (& resolution).

# Our motivation and high-level overview of results

---

**Motivation:** Construct cohomogeneity one Laplacian solitons that are potentially viable models for finite-time singularity formation (& resolution).

**Goal:** Find **complete**  $G_2$  solitons with **cohomogeneity one**, specifically

- $SU(3)$ -invariant ones on  $\Lambda_-^2 \mathbb{C}P^2$     &    •  $Sp(2)$ -invariant ones on  $\Lambda_-^2 \mathbb{S}^4$
- and understand their asymptotic geometry.

# Our motivation and high-level overview of results

---

**Motivation:** Construct cohomogeneity one Laplacian solitons that are potentially viable models for finite-time singularity formation (& resolution).

**Goal:** Find **complete**  $G_2$  solitons with **cohomogeneity one**, specifically

- $SU(3)$ -invariant ones on  $\Lambda_-^2 \mathbb{C}P^2$  & •  $Sp(2)$ -invariant ones on  $\Lambda_-^2 \mathbb{S}^4$

and understand their asymptotic geometry.

## Telegraphic overview of results:

- We find complete shrinkers, steady solitons and expanders
- Shrinkers are the rarest and most rigid
- Expanders are the most abundant
- Steady solitons are intermediate between shrinkers and expanders; our steady solitons behave very differently to steady Ricci solitons.

# Our motivation and high-level overview of results

---

**Motivation:** Construct cohomogeneity one Laplacian solitons that are potentially viable models for finite-time singularity formation (& resolution).

**Goal:** Find **complete**  $G_2$  solitons with **cohomogeneity one**, specifically

- $SU(3)$ -invariant ones on  $\Lambda_-^2 \mathbb{C}P^2$  & •  $Sp(2)$ -invariant ones on  $\Lambda_-^2 \mathbb{S}^4$

and understand their asymptotic geometry.

## Telegraphic overview of results:

- We find complete shrinkers, steady solitons and expanders
- Shrinkers are the rarest and most rigid
- Expanders are the most abundant
- Steady solitons are intermediate between shrinkers and expanders; our steady solitons behave very differently to steady Ricci solitons.
- Almost all our complete examples are **asymptotically conical (AC)**; **but** there are complete steady solitons with **exponential volume growth** that appear at the boundary of the space of AC steady solitons. Related to scalar curvature of closed  $G_2$ -structures being non-positive.

# Our motivation and high-level overview of results

---

**Motivation:** Construct cohomogeneity one Laplacian solitons that are potentially viable models for finite-time singularity formation (& resolution).

**Goal:** Find **complete**  $G_2$  solitons with **cohomogeneity one**, specifically

- $SU(3)$ -invariant ones on  $\Lambda_-^2 \mathbb{C}P^2$  & •  $Sp(2)$ -invariant ones on  $\Lambda_-^2 \mathbb{S}^4$

and understand their asymptotic geometry.

## Telegraphic overview of results:

- We find complete shrinkers, steady solitons and expanders
- Shrinkers are the rarest and most rigid
- Expanders are the most abundant
- Steady solitons are intermediate between shrinkers and expanders; our steady solitons behave very differently to steady Ricci solitons.
- Almost all our complete examples are **asymptotically conical (AC)**; **but** there are complete steady solitons with **exponential volume growth** that appear at the boundary of the space of AC steady solitons.
- There are important differences between the  $Sp(2)$  and  $SU(3)$ -invariant cases.

# Theorems and conjectures on Laplacian shrinkers

---

## Theorem Shrink1

There exists an explicit complete AC shrinker with rate  $-1$  on  $\Lambda^2 \mathbb{S}^4$  and on  $\Lambda^2 \mathbb{C}P^2$ .

- Shrinkers are rare! Possible models for *formation of conical singularities*.

# Theorems and conjectures on Laplacian shrinkers

---

## Theorem Shrink1

There exists an explicit complete AC shrinker with rate  $-1$  on  $\Lambda^2 \mathbb{S}^4$  and on  $\Lambda^2 \mathbb{C}P^2$ .

- Shrinkers are rare! Possible models for *formation of conical singularities*.

## Theorem Shrink2

Let  $G$  be  $SU(3)$  or  $Sp(2)$ . For every closed  $G$ -invariant  $G_2$ -cone  $C$  there exists a unique  $G$ -invariant shrinker AC end (i.e. need not extend to a complete AC shrinker) asymptotic to  $C$ .

- The space of  $G$ -invariant  $G_2$ -cones is 1-dimensional for  $G = Sp(2)$  and 2-dimensional for  $G = SU(3)$ . So previous theorem yields continuous families of AC shrinker *ends*.

# Theorems and conjectures on Laplacian shrinkers

---

## Theorem Shrink1

There exists an explicit complete AC shrinker with rate  $-1$  on  $\Lambda^2 \mathbb{S}^4$  and on  $\Lambda^2 \mathbb{C}P^2$ .

- Shrinkers are rare! Possible models for *formation of conical singularities*.

## Theorem Shrink2

Let  $G$  be  $SU(3)$  or  $Sp(2)$ . For every closed  $G$ -invariant  $G_2$ -cone  $C$  there exists a unique  $G$ -invariant shrinker AC end (i.e. need not extend to a complete AC shrinker) asymptotic to  $C$ .

- The space of  $G$ -invariant  $G_2$ -cones is 1-dimensional for  $G = Sp(2)$  and 2-dimensional for  $G = SU(3)$ . So previous theorem yields continuous families of AC shrinker ends.

## Shrinker Conjectures

- (i) The explicit  $Sp(2)$ -invariant AC shrinker on  $\Lambda^2 \mathbb{S}^4$  is the unique complete AC  $Sp(2)$ -invariant shrinker.
- (ii) The explicit  $Sp(2)$ -invariant AC shrinker on  $\Lambda^2 \mathbb{S}^4$  is the unique complete AC shrinker asymptotic to an  $Sp(2)$ -invariant closed  $G_2$ -cone, i.e. we assume symmetry only of the cone not of the AC shrinker.
- (iii) There are only finitely many complete  $SU(3)$ -invariant AC shrinkers.

So MOST AC shrinker ends from **Thm Shrink2** SHOULDN'T extend to complete AC shrinkers.



# Theorems and conjectures on Laplacian expanders

---

## Theorem Expand1

- (i) There exists a 1-parameter family of complete  $Sp(2)$ -invariant AC expanders with rate  $-1$  on  $\Lambda^2 \mathbb{S}^4$  and a 1-parameter family of  $SU(3) \times \mathbb{Z}_2$ -invariant AC expanders on  $\Lambda^2 \mathbb{C}P^2$ .
- (ii) Every  $Sp(2)$ -invariant closed  $G_2$ -cone 'on one side' of the torsion-free cone arises as the asymptotic cone of a unique (up to scale) complete AC  $Sp(2)$ -invariant expander.
- AC expanders give models for how Laplacian flow can *smooth out certain conical singularities*.

# Theorems and conjectures on Laplacian expanders

---

## Theorem Expand1

(i) There exists a 1-parameter family of complete  $Sp(2)$ -invariant AC expanders with rate  $-1$  on  $\Lambda^2 \mathbb{S}^4$  and a 1-parameter family of  $SU(3) \times \mathbb{Z}_2$ -invariant AC expanders on  $\Lambda^2 \mathbb{C}P^2$ .

(ii) Every  $Sp(2)$ -invariant closed  $G_2$ -cone 'on one side' of the torsion-free cone arises as the asymptotic cone of a unique (up to scale) complete AC  $Sp(2)$ -invariant expander.

- AC expanders give models for how Laplacian flow can *smooth out certain conical singularities*.

## Theorem Expand2

Let  $G$  be  $SU(3)$  or  $Sp(2)$  and  $k$  be the dimension of the space of  $G$ -invariant closed  $G_2$ -cones.

$\exists$  a  $k$ -diml family of  $G$ -invariant AC expander ends asymptotic to any closed  $G$ -invariant  $G_2$ -cone; the difference between two such AC expanders is of order  $\exp(-\frac{\lambda}{6} t^2) \times$  polynomial, where  $\lambda > 0$  is the dilation constant of the expander.

# Theorems and conjectures on Laplacian expanders

---

## Theorem Expand1

(i) There exists a 1-parameter family of complete  $Sp(2)$ -invariant AC expanders with rate  $-1$  on  $\Lambda^2 \mathbb{S}^4$  and a 1-parameter family of  $SU(3) \times \mathbb{Z}_2$ -invariant AC expanders on  $\Lambda^2 \mathbb{C}P^2$ .

(ii) Every  $Sp(2)$ -invariant closed  $G_2$ -cone 'on one side' of the torsion-free cone arises as the asymptotic cone of a unique (up to scale) complete AC  $Sp(2)$ -invariant expander.

• AC expanders give models for how Laplacian flow can *smooth out certain conical singularities*.

## Theorem Expand2

Let  $G$  be  $SU(3)$  or  $Sp(2)$  and  $k$  be the dimension of the space of  $G$ -invariant closed  $G_2$ -cones.

$\exists$  a  $k$ -diml family of  $G$ -invariant AC expander ends asymptotic to any closed  $G$ -invariant  $G_2$ -cone; the difference between two such AC expanders is of order  $\exp(-\frac{\lambda}{6} t^2) \times$  polynomial, where  $\lambda > 0$  is the dilation constant of the expander.

## Expander Conjectures

(i) There is a 2-parameter family of complete  $SU(3)$ -invariant AC expanders on  $\Lambda^2 \mathbb{C}P^2$ .

(ii) The set of asymptotic cones of complete  $SU(3)$ -invariant AC expanders is a proper open subset of the 2-diml space of all  $SU(3)$ -invariant closed  $G_2$ -cones.

We have a precise conjecture for what this proper open subset should be.

## $Sp(2)$ -invariant singularity formation

In other geometric flows, often there exists an AC shrinker and an AC expander that share a common (asymptotic) cone; combining them yields a 'weak solution' to the flow that is singular only at the time instant  $t = 0$  where the common cone appears.

Feldman-Ilmanen-Knopf (FIK) called this 'flowing through the singularity'.

e.g. in Kähler-Ricci flow there is such an AC shrinker (due to FIK) / AC expander (due to Cao) pair for blowing-down a  $(-1)$  curve in a Kähler surface

## $Sp(2)$ -invariant singularity formation

In other geometric flows, often there exists an AC shrinker and an AC expander that share a common (asymptotic) cone; combining them yields a 'weak solution' to the flow that is singular only at the time instant  $t = 0$  where the common cone appears.

Feldman-Ilmanen-Knopf (FIK) called this 'flowing through the singularity'.

### **No-Flow-Through Theorem**

There is no complete  $Sp(2)$ -invariant AC expander whose cone coincides with the cone of the explicit AC  $Sp(2)$ -invariant shrinker on  $\Lambda_-^2 \mathbb{S}^4$ .

*We conjecture the result holds without the  $Sp(2)$ -invariance assumption on the AC expander.*

## $Sp(2)$ -invariant singularity formation

In other geometric flows, often there exists an AC shrinker and an AC expander that share a common (asymptotic) cone; combining them yields a 'weak solution' to the flow that is singular only at the time instant  $t = 0$  where the common cone appears.

Feldman-Ilmanen-Knopf (FIK) called this 'flowing through the singularity'.

### No-Flow-Through Theorem

There is no complete  $Sp(2)$ -invariant AC expander whose cone coincides with the cone of the explicit AC  $Sp(2)$ -invariant shrinker on  $\Lambda_-^2 \mathbb{S}^4$ .

*We conjecture the result holds without the  $Sp(2)$ -invariance assumption on the AC expander.*

### Obvious questions this raises:

1. Does the explicit AC  $Sp(2)$ -invariant shrinker arise as the singularity model for a finite-time singularity of Laplacian flow for some closed  $G_2$ -structure on a compact 7-manifold?
2. If **yes** to **Q1** how should we continue Laplacian flow past this singularity?

Should we consider a singular version of Laplacian flow where the conical singularity persists but can vary within the set of closed  $G_2$ -cones?

How does the bulk geometry drive the evolution of the conical singularity?

## $Sp(2)$ -invariant singularity formation

In other geometric flows, often there exists an AC shrinker and an AC expander that share a common (asymptotic) cone; combining them yields a 'weak solution' to the flow that is singular only at the time instant  $t = 0$  where the common cone appears.

Feldman-Ilmanen-Knopf (FIK) called this 'flowing through the singularity'.

### No-Flow-Through Theorem

There is no complete  $Sp(2)$ -invariant AC expander whose cone coincides with the cone of the explicit AC  $Sp(2)$ -invariant shrinker on  $\Lambda_-^2 \mathbb{S}^4$ .

*We conjecture the result holds without the  $Sp(2)$ -invariance assumption on the AC expander.*

### Obvious questions this raises:

1. Does the explicit AC  $Sp(2)$ -invariant shrinker arise as the singularity model for a finite-time singularity of Laplacian flow for some closed  $G_2$ -structure on a compact 7-manifold?
2. If **yes** to **Q1** how should we continue Laplacian flow past this singularity?

Should we consider a singular version of Laplacian flow where the conical singularity persists but can vary within the set of closed  $G_2$ -cones?

How does the bulk geometry drive the evolution of the conical singularity?

- We expect *different* behaviour for the explicit  $SU(3)$ -invariant AC shrinker on  $\Lambda_-^2 \mathbb{C}P^2$ !

# $SU(3)$ -invariant steady Laplacian solitons

## Theorem Steady

There exists (up to scale) a 1-parameter family of complete  $SU(3)$ -invariant steady solitons on  $\Lambda_-^2 \mathbb{C}P^2$  parameterised by  $s \in [-1, 1]$ .

- (i) For  $s = 0$  it is the standard  $SU(3)$ -invariant torsion-free AC  $G_2$  structure on  $\Lambda_-^2 \mathbb{C}P^2$ .
- (ii) For  $\pm s \in (0, 1)$  it is a nontrivial steady soliton asymptotic with rate  $-1$  to the unique  $SU(3)$ -invariant torsion-free cone.
- (iii) For  $s = \pm 1$  it is an *explicit* nontrivial steady soliton with *exponential volume growth*.

Asymptotically this steady soliton has constant negative scalar curvature and approaches a flat  $T^2$ -bundle over the sinh-cone over  $(\mathbb{C}P^2, g_{FS})$  where  $g_{FS}$  denotes the Fubini-Study metric.

(The sinh-cone of  $\mathbb{C}P^2$  is a noncompact Einstein 5-manifold with negative Einstein constant.)



# $SU(3)$ -invariant steady Laplacian solitons

## Theorem Steady

There exists (up to scale) a 1-parameter family of complete  $SU(3)$ -invariant steady solitons on  $\Lambda^2 \mathbb{C}P^2$  parameterised by  $s \in [-1, 1]$ .

- (i) For  $s = 0$  it is the standard  $SU(3)$ -invariant torsion-free AC  $G_2$  structure on  $\Lambda^2 \mathbb{C}P^2$ .
- (ii) For  $\pm s \in (0, 1)$  it is a nontrivial steady soliton asymptotic with rate  $-1$  to the unique  $SU(3)$ -invariant torsion-free cone.
- (iii) For  $s = \pm 1$  it is an *explicit* nontrivial steady soliton with *exponential volume growth*.

Asymptotically this steady soliton has constant negative scalar curvature and approaches a flat  $T^2$ -bundle over the sinh-cone over  $(\mathbb{C}P^2, g_{FS})$  where  $g_{FS}$  denotes the Fubini-Study metric.

## Remarks:

- Any complete  $Sp(2)$ -invariant steady soliton must be a trivial steady soliton.
- **AC** steady solitons a new feature compared to Ricci/Kähler-Ricci flow.
- the asymptotic behaviour of the explicit steady soliton is impossible for steady Ricci solitons because they have scalar curvature  $S \geq 0$ .
- Understanding complete  $SU(3)$ -invariant steady solitons is also useful for proving results about complete  $SU(3)$ -invariant expanders, e.g. via blow-down arguments.

## Closed invariant $G_2$ -structures on $\Lambda_-^2 M^4 \setminus M$

---

For  $M = \mathbb{C}P^2$  or  $\mathbb{S}^4$ :

- $\Lambda_-^2 M$  has a cohomogeneity one action by  $G = SU(3)$  or  $Sp(2)$
- $\Lambda_-^2 M \setminus M$  is diffeomorphic to  $\mathbb{R}_+ \times \Sigma$ , for  $\Sigma = SU(3)/T^2$  or  $\mathbb{C}P^3$

## Closed invariant $G_2$ -structures on $\Lambda_-^2 M^4 \setminus M$

---

For  $M = \mathbb{C}P^2$  or  $\mathbb{S}^4$ :

- $\Lambda_-^2 M$  has a cohomogeneity one action by  $G = SU(3)$  or  $Sp(2)$
- $\Lambda_-^2 M \setminus M$  is diffeomorphic to  $\mathbb{R}_+ \times \Sigma$ , for  $\Sigma = SU(3)/T^2$  or  $\mathbb{C}P^3$

Any closed  $G$ -invariant  $G_2$ -structure  $\varphi$  on  $\mathbb{R}_+ \times \Sigma$  can be written as

$$\varphi = (f_1^2 \omega_1 + f_2^2 \omega_2 + f_3^2 \omega_3) \wedge dt + f_1 f_2 f_3 \alpha = \omega_f \wedge dt + \operatorname{Re} \Omega_f$$

where  $\omega_1, \omega_2, \omega_3 \in \Omega^2(\Sigma)$  and  $\alpha \in \Omega^3(\Sigma)$  are  $G$ -invariant forms on  $\Sigma$  and  $f_i : \mathbb{R}_+ \rightarrow \mathbb{R}_+$  satisfy

$$(f_1 f_2 f_3)' = \frac{1}{2}(f_1^2 + f_2^2 + f_3^2) \quad (\#)$$

## Closed invariant $G_2$ -structures on $\Lambda_-^2 M^4 \setminus M$

---

For  $M = \mathbb{C}P^2$  or  $\mathbb{S}^4$ :

- $\Lambda_-^2 M$  has a cohomogeneity one action by  $G = SU(3)$  or  $Sp(2)$
- $\Lambda_-^2 M \setminus M$  is diffeomorphic to  $\mathbb{R}_+ \times \Sigma$ , for  $\Sigma = SU(3)/T^2$  or  $\mathbb{C}P^3$

Any closed  $G$ -invariant  $G_2$ -structure  $\varphi$  on  $\mathbb{R}_+ \times \Sigma$  can be written as

$$\varphi = (f_1^2 \omega_1 + f_2^2 \omega_2 + f_3^2 \omega_3) \wedge dt + f_1 f_2 f_3 \alpha = \omega_f \wedge dt + \operatorname{Re} \Omega_f$$

where  $\omega_1, \omega_2, \omega_3 \in \Omega^2(\Sigma)$  and  $\alpha \in \Omega^3(\Sigma)$  are  $G$ -invariant forms on  $\Sigma$  and  $f_i : \mathbb{R}_+ \rightarrow \mathbb{R}_+$  satisfy

$$(f_1 f_2 f_3)' = \frac{1}{2}(f_1^2 + f_2^2 + f_3^2) \quad (\#)$$

For  $Sp(2)$ -invariance in addition we require  $f_2 = f_3$ . Structure equations for  $\omega_i, \alpha$  the same in both cases  $\Rightarrow \Lambda_-^2 \mathbb{S}^4$  case can be treated as a special case of  $\Lambda_-^2 \mathbb{C}P^2$  case where  $f_2 = f_3$ .

## Closed invariant $G_2$ -structures on $\Lambda_-^2 M^4 \setminus M$

---

For  $M = \mathbb{C}P^2$  or  $\mathbb{S}^4$ :

- $\Lambda_-^2 M$  has a cohomogeneity one action by  $G = SU(3)$  or  $Sp(2)$
- $\Lambda_-^2 M \setminus M$  is diffeomorphic to  $\mathbb{R}_+ \times \Sigma$ , for  $\Sigma = SU(3)/T^2$  or  $\mathbb{C}P^3$

Any closed  $G$ -invariant  $G_2$ -structure  $\varphi$  on  $\mathbb{R}_+ \times \Sigma$  can be written as

$$\varphi = (f_1^2 \omega_1 + f_2^2 \omega_2 + f_3^2 \omega_3) \wedge dt + f_1 f_2 f_3 \alpha = \omega_f \wedge dt + \operatorname{Re} \Omega_f$$

where  $\omega_1, \omega_2, \omega_3 \in \Omega^2(\Sigma)$  and  $\alpha \in \Omega^3(\Sigma)$  are  $G$ -invariant forms on  $\Sigma$  and  $f_i : \mathbb{R}_+ \rightarrow \mathbb{R}_+$  satisfy

$$(f_1 f_2 f_3)' = \frac{1}{2}(f_1^2 + f_2^2 + f_3^2) \quad (\#)$$

For  $Sp(2)$ -invariance in addition we require  $f_2 = f_3$ . Structure equations for  $\omega_i, \alpha$  the same in both cases  $\Rightarrow \Lambda_-^2 \mathbb{S}^4$  case can be treated as a special case of  $\Lambda_-^2 \mathbb{C}P^2$  case where  $f_2 = f_3$ .

*Discrete symmetries for  $G = SU(3)$ :* When  $f_j = f_k$  for  $j \neq k$  metric on the corresponding principal orbit has an extra free isometric  $\mathbb{Z}_2$ -action, which does NOT preserve the  $SU(3)$  structure  $(\omega_f, \Omega_f)$ . When  $f_1 = f_2 = f_3$  the orbit has an additional free isometric action of the symmetric group  $S_3$ , and the subgroup  $A_3 < S_3$  preserves  $(\omega_f, \Omega_f)$ .

## Closed invariant $G_2$ -structures on $\Lambda_-^2 M^4 \setminus M$

---

For  $M = \mathbb{C}P^2$  or  $\mathbb{S}^4$ :

- $\Lambda_-^2 M$  has a cohomogeneity one action by  $G = SU(3)$  or  $Sp(2)$
- $\Lambda_-^2 M \setminus M$  is diffeomorphic to  $\mathbb{R}_+ \times \Sigma$ , for  $\Sigma = SU(3)/T^2$  or  $\mathbb{C}P^3$

Any closed  $G$ -invariant  $G_2$ -structure  $\varphi$  on  $\mathbb{R}_+ \times \Sigma$  can be written as

$$\varphi = (f_1^2 \omega_1 + f_2^2 \omega_2 + f_3^2 \omega_3) \wedge dt + f_1 f_2 f_3 \alpha = \omega_f \wedge dt + \operatorname{Re} \Omega_f$$

where  $\omega_1, \omega_2, \omega_3 \in \Omega^2(\Sigma)$  and  $\alpha \in \Omega^3(\Sigma)$  are  $G$ -invariant forms on  $\Sigma$  and  $f_i : \mathbb{R}_+ \rightarrow \mathbb{R}_+$  satisfy

$$(f_1 f_2 f_3)' = \frac{1}{2}(f_1^2 + f_2^2 + f_3^2) \quad (\#)$$

For  $Sp(2)$ -invariance in addition we require  $f_2 = f_3$ . Structure equations for  $\omega_i, \alpha$  the same in both cases  $\Rightarrow \Lambda_-^2 \mathbb{S}^4$  case can be treated as a special case of  $\Lambda_-^2 \mathbb{C}P^2$  case where  $f_2 = f_3$ .

*Discrete symmetries for  $G = SU(3)$ :* When  $f_j = f_k$  for  $j \neq k$  metric on the corresponding principal orbit has an extra free isometric  $\mathbb{Z}_2$ -action, which does NOT preserve the  $SU(3)$  structure  $(\omega_f, \Omega_f)$ . When  $f_1 = f_2 = f_3$  the orbit has an additional free isometric action of the symmetric group  $S_3$ , and the subgroup  $A_3 < S_3$  preserves  $(\omega_f, \Omega_f)$ .

$$(\#) \Rightarrow \frac{d}{dt}(f_1 f_2 f_3)^{1/3} \geq \frac{1}{2} \quad \text{i.e. volume grows at least like } t^7.$$

## The soliton equations and their local solutions

---

The **torsion**  $\tau$  of a closed  $G_2$ -structure  $\varphi$  is the 2-form (of type 14) s.t.

$$d(*\varphi) = \tau \wedge \varphi.$$

In our  $G$ -invariant setting the torsion of  $\varphi_f$  is  $\tau = \sum \tau_i \omega_i$  where  $\omega_i$  are the  $G$ -invariant 2-forms on  $\Sigma$  and

$$\tau_i = (f_i^2)' + \frac{f_i^2}{f_1 f_2 f_3} \left( 2f_i^2 - \sum f_i^2 \right). \quad (1)$$

# The soliton equations and their local solutions

---

The **torsion**  $\tau$  of a closed  $G_2$ -structure  $\varphi$  is the 2-form (of type 14) s.t.

$$d(*\varphi) = \tau \wedge \varphi.$$

In our  $G$ -invariant setting the torsion of  $\varphi_f$  is  $\tau = \sum \tau_i \omega_i$  where  $\omega_i$  are the  $G$ -invariant 2-forms on  $\Sigma$  and

$$\tau_i = (f_i^2)' + \frac{f_i^2}{f_1 f_2 f_3} \left( 2f_i^2 - \sum f_j^2 \right). \quad (1)$$

The soliton condition for  $(\varphi_f, X = u \frac{\partial}{\partial t}, \lambda)$  is the (mixed-order) ODE system

$$2(f_1 f_2 f_3)' = f_1^2 + f_2^2 + f_3^2, \quad (2a)$$

$$(\tau_i - u f_i^2)' = \lambda f_i^2, \quad \text{for } i = 1, 2, 3, \quad (2b)$$

$$\tau_1 + \tau_2 + \tau_3 = u(f_1^2 + f_2^2 + f_3^2) + 2\lambda f_1 f_2 f_3. \quad (2c)$$



# The soliton equations and their local solutions

---

The **torsion**  $\tau$  of a closed  $G_2$ -structure  $\varphi$  is the 2-form (of type 14) s.t.

$$d(*\varphi) = \tau \wedge \varphi.$$

In our  $G$ -invariant setting the torsion of  $\varphi_f$  is  $\tau = \sum \tau_i \omega_i$  where  $\omega_i$  are the  $G$ -invariant 2-forms on  $\Sigma$  and

$$\tau_i = (f_i^2)' + \frac{f_i^2}{f_1 f_2 f_3} \left( 2f_i^2 - \sum f_j^2 \right). \quad (1)$$

The soliton condition for  $(\varphi_f, X = u \frac{\partial}{\partial t}, \lambda)$  is the (mixed-order) ODE system

$$2(f_1 f_2 f_3)' = f_1^2 + f_2^2 + f_3^2, \quad (2a)$$

$$(\tau_i - u f_i^2)' = \lambda f_i^2, \quad \text{for } i = 1, 2, 3, \quad (2b)$$

$$\tau_1 + \tau_2 + \tau_3 = u(f_1^2 + f_2^2 + f_3^2) + 2\lambda f_1 f_2 f_3. \quad (2c)$$

We can rewrite (2) as a real analytic *1st-order system in the 6 variables*  $(f_i, \tau_i)$  and the type 14 condition on  $\tau$  imposes the algebraic condition  $\sum \tau_i f_j^2 f_k^2 = 0$ .

Then  $u$  is determined algebraically from  $(f, \tau)$  by equation (2c).

$\Rightarrow \exists$  a 4-diml family of  $SU(3)$ -invariant solitons & 2-diml family of  $Sp_2$ -invariant solitons.

In the *steady case*  $\lambda = 0$  the action of scaling reduces these parameter counts by 1.

## Smooth extension over the zero section of $\Lambda_-^2 M^4$

---

Understand solutions defined near zero section of  $\Lambda_-^2 M$  that extend smoothly over it.

### Proposition

- For each  $\lambda \in \mathbb{R}$ , there is a 2-parameter family  $\varphi_{b,c}$  of solutions defined for small  $t$  that extend smoothly to a  $\lambda$ -soliton on (nhd of zero section in)  $\Lambda_-^2 \mathbb{C}P^2$ ;
- the 1-parameter subfamily  $\varphi_b = \varphi_{b,0}$  also defines  $\lambda$ -solitons on  $\Lambda_-^2 \mathbb{S}^4$ .

Two scale-invariant parameters:  $\lambda b^2$  and  $c$ . So up to scale:

- 2-parameter families of smoothly-closing expanders/shrinkers on  $\Lambda_-^2 \mathbb{C}P^2$
- a 1-parameter family of smoothly-closing steady solitons on  $\Lambda_-^2 \mathbb{C}P^2$

## Smooth extension over the zero section of $\Lambda_-^2 M^4$

---

Understand solutions defined near zero section of  $\Lambda_-^2 M$  that extend smoothly over it.

### Proposition

- For each  $\lambda \in \mathbb{R}$ , there is a 2-parameter family  $\varphi_{b,c}$  of solutions defined for small  $t$  that extend smoothly to a  $\lambda$ -soliton on (nhd of zero section in)  $\Lambda_-^2 \mathbb{C}P^2$ ;
- the 1-parameter subfamily  $\varphi_b = \varphi_{b,0}$  also defines  $\lambda$ -solitons on  $\Lambda_-^2 \mathbb{S}^4$ .

Two scale-invariant parameters:  $\lambda b^2$  and  $c$ . So up to scale:

- 2-parameter families of smoothly-closing expanders/shrinkers on  $\Lambda_-^2 \mathbb{C}P^2$
- a 1-parameter family of smoothly-closing steady solitons on  $\Lambda_-^2 \mathbb{C}P^2$

Idea of proof: The closed  $G$ -invariant  $G_2$ -structure

$$\varphi = (f_1^2 \omega_1 + f_2^2 \omega_2 + f_3^2 \omega_3) \wedge dt + f_1 f_2 f_3 \alpha$$

on  $\mathbb{R}_+ \times \Sigma$  extends to smooth  $G_2$ -structure on  $\Lambda_-^2 M$  if and only if

$f_1$  is odd with  $f_1'(0) = 1$ , and  $f_2$  and  $f_3$  are even with  $b := f_2(0) = f_3(0) \neq 0$ .

Resulting singular initial value problem for the higher order terms has formal power series solutions that are convergent. (It is a **regular singular point** of 1st-order ODE system).

# Proof sketch of Theorem Steady

---

## 1. Decoupling

- For  $\lambda = 0$ , the ODEs can be separated into evolution of *scale*  $g$  and evolution of 4 scale-normalised variables.
- Unique fixed point for scale-normalised flow is torsion-free cone; it is a *stable* fixed point.

## 2. Smoothly-closing solutions

- Near special orbit  $\mathbb{C}P^2$ ,  $\exists$  a 1-parameter family of solutions  $\varphi_c$  up to scale.
- Unique one with  $f_2 = f_3$ : static soliton from Bryant–Salamon AC  $G_2$ -mfd has  $c = 0$ .

1 & 2: Stability of fixed point and continuous dependence of smoothly-closing solutions on  $c$   
 $\Rightarrow$  persistence of AC asymptotics for  $c$  sufficiently small.

## 3. The explicit solution

Numerical simulations suggested critical value  $c_{crit}$  of  $c$  at which AC asymptotics terminated. Inspection of power series solutions for  $c_{crit}$  led to initial guess for explicit solution.

## 4. Trapping by the explicit solution

Evolution of a quantity  $G$  suggested by the explicit solution ( $G$  is constant on it) guarantees that for any  $c < c_{crit}$  the smoothly-closing solution  $\varphi_c$  is complete and has AC asymptotics.

# Shrinkers: consequences of AC end rigidity (Thm Shrink2)

---

**Heuristic for  $\lambda < 0$ :** Invariant shrinkers on  $\mathbb{R}_+ \times SU(3)/T^2$  are flow lines in 5-dim phase space.

In 4-dimensional space of flow lines

- Proposition  $\Rightarrow$  a 2-dimensional submanifold of solutions extend across zero section

$$\mathbb{C}P^2 \subset \Lambda_-^2 \mathbb{C}P^2$$

- Thm Shrink2  $\Rightarrow$  a 2-dimensional submanifold of solutions has AC behaviour

Expect transverse intersections  $\rightsquigarrow$  *finitely* many AC shrinkers on  $\Lambda_-^2 \mathbb{C}P^2$ .

# Shrinkers: consequences of AC end rigidity (Thm Shrink2)

---

**Heuristic for  $\lambda < 0$ :** Invariant shrinkers on  $\mathbb{R}_+ \times SU(3)/T^2$  are flow lines in 5-dim phase space.

In 4-dimensional space of flow lines

- Proposition  $\Rightarrow$  a 2-dimensional submanifold of solutions extend across zero section

$$\mathbb{C}P^2 \subset \Lambda_-^2 \mathbb{C}P^2$$

- Thm Shrink2  $\Rightarrow$  a 2-dimensional submanifold of solutions has AC behaviour

Expect transverse intersections  $\rightsquigarrow$  *finitely* many AC shrinkers on  $\Lambda_-^2 \mathbb{C}P^2$ .

Similarly, restricting attention to solutions with  $f_2 = f_3$ :

- 2-dimensional space of flow lines;
- 1-dim submanifold extends over special orbit; 1-dim submanifold has AC behaviour.

Expect transverse intersections  $\rightsquigarrow$  *finitely* many AC shrinkers on  $\Lambda_-^2 \mathbb{S}^4$ .

# Shrinkers: consequences of AC end rigidity (Thm Shrink2)

**Heuristic for  $\lambda < 0$ :** Invariant shrinkers on  $\mathbb{R}_+ \times SU(3)/T^2$  are flow lines in 5-dim phase space.

In 4-dimensional space of flow lines

- Proposition  $\Rightarrow$  a 2-dimensional submanifold of solutions extend across zero section

$$\mathbb{C}P^2 \subset \Lambda_-^2 \mathbb{C}P^2$$

- Thm Shrink2  $\Rightarrow$  a 2-dimensional submanifold of solutions has AC behaviour

Expect transverse intersections  $\rightsquigarrow$  *finitely* many AC shrinkers on  $\Lambda_-^2 \mathbb{C}P^2$ .

Similarly, restricting attention to solutions with  $f_2 = f_3$ :

- 2-dimensional space of flow lines;
- 1-dim submanifold extends over special orbit; 1-dim submanifold has AC behaviour.

Expect transverse intersections  $\rightsquigarrow$  *finitely* many AC shrinkers on  $\Lambda_-^2 \mathbb{S}^4$ .

In fact, can spot one explicit solution!

**Theorem Shrink1:** For  $\lambda = -1$

$$f_1 = t, \quad f_2^2 = f_3^2 = \frac{9}{4} + \frac{1}{4}t^2, \quad u = \frac{t}{3} + \frac{4t}{9 + t^2}$$

is an AC shrinker with rate  $-1$  asymptotic to the cone  $(1, \frac{1}{2}, \frac{1}{2})$ .

# Shrinkers: consequences of AC end rigidity (Thm Shrink2)

**Heuristic for  $\lambda < 0$ :** Invariant shrinkers on  $\mathbb{R}_+ \times SU(3)/T^2$  are flow lines in 5-dim phase space.

In 4-dimensional space of flow lines

- Proposition  $\Rightarrow$  a 2-dimensional submanifold of solutions extend across zero section

$$\mathbb{C}P^2 \subset \Lambda_-^2 \mathbb{C}P^2$$

- Thm Shrink2  $\Rightarrow$  a 2-dimensional submanifold of solutions has AC behaviour

Expect transverse intersections  $\rightsquigarrow$  *finitely* many AC shrinkers on  $\Lambda_-^2 \mathbb{C}P^2$ .

Similarly, restricting attention to solutions with  $f_2 = f_3$ :

- 2-dimensional space of flow lines;
- 1-dim submanifold extends over special orbit; 1-dim submanifold has AC behaviour.

Expect transverse intersections  $\rightsquigarrow$  *finitely* many AC shrinkers on  $\Lambda_-^2 \mathbb{S}^4$ .

In fact, can spot one explicit solution!

**Theorem Shrink1:** For  $\lambda = -1$

$$f_1 = t, \quad f_2^2 = f_3^2 = \frac{9}{4} + \frac{1}{4}t^2, \quad u = \frac{t}{3} + \frac{4t}{9 + t^2}$$

is an AC shrinker with rate  $-1$  asymptotic to the cone  $(1, \frac{1}{2}, \frac{1}{2})$ .

**Conjecture:** *This is the unique  $Sp_2$ -invariant AC shrinker on  $\Lambda_-^2 \mathbb{S}^4$ .*



**Thanks for your attention!**

**Bonus slides!**

# Expectations about $SU(3)$ -invariant expanders.

---

## Conjecture

(i) There is a 2-parameter family of complete  $SU(3)$ -invariant expanders on  $\Lambda_-^2 \mathbb{C}P^2$  asymptotic to a closed non-torsion free  $SU(3)$ -invariant  $G_2$ -cone.

(ii) The set of possible asymptotic cones of complete  $SU(3)$ -invariant expanders on  $\Lambda_-^2 \mathbb{C}P^2$  (with singular orbit type  $\mathbb{C}P_1^2$ , i.e.  $f_1 \rightarrow 0$  as  $t \rightarrow 0$ ) is a proper connected open subset of the 2-dimensional space of all closed  $SU(3)$ -invariant  $G_2$ -cones; it is bounded by those cones with  $f_1 = f_2$  or  $f_1 = f_3$ , i.e. cones that possess either of the two possible extra  $\mathbb{Z}_2$ -isometries that differ from the extra asymptotic  $\mathbb{Z}_2$ -isometry present close to the singular orbit  $\mathbb{C}P_1^2$ .

## Corollary

There are two distinct complete AC  $SU(3)$ -invariant expanders on  $\Lambda_-^2 \mathbb{C}P^2$  asymptotic to the cone of the explicit  $SU(3)$ -invariant shrinker on  $\Lambda_-^2 \mathbb{C}P^2$  constructed in Theorem Shrink1. They differ by their singular orbit types:  $\mathbb{C}P_2^2$  and  $\mathbb{C}P_3^2$  (versus  $\mathbb{C}P_1^2$ ).

$\Rightarrow$  In the  $SU(3)$ -invariant case we CAN 'flow through the singularity': there is a 'weak solution' to Laplacian flow obtained by using the explicit  $SU(3)$ -invariant AC shrinker for  $t < 0$ , the cone for  $t = 0$  and either of the two compatible AC expanders for  $t > 0$ .

We clearly don't have a unique weak solution in this case.

## Constructing invariant non-steady AC ends

---

The problem of constructing an *invariant non-steady soliton AC end asymptotic to a given closed invariant cone* can be written as another singular initial value problem (SIVP) for the first-order ODE system: this time the SIVP is **irregular**.

## Constructing invariant non-steady AC ends

---

The problem of constructing an *invariant non-steady soliton AC end asymptotic to a given closed invariant cone* can be written as another singular initial value problem (SIVP) for the first-order ODE system: this time the SIVP is **irregular**.

**Theorem 1** For any  $\lambda \neq 0$ ,  $\exists$  a ! formal power series solution  $\mathcal{P}$  in  $t^{-1}$  determined by the cone;  $\exists$  a solution of the ODE system that is smooth in a nhd of  $t = +\infty$  and whose Taylor series is  $\mathcal{P}$ .

## Constructing invariant non-steady AC ends

---

The problem of constructing an *invariant non-steady soliton AC end asymptotic to a given closed invariant cone* can be written as another singular initial value problem (SIVP) for the first-order ODE system: this time the SIVP is **irregular**.

**Theorem 1** For any  $\lambda \neq 0$ ,  $\exists$  a ! formal power series solution  $\mathcal{P}$  in  $t^{-1}$  determined by the cone;  $\exists$  a solution of the ODE system that is smooth in a nhd of  $t = +\infty$  and whose Taylor series is  $\mathcal{P}$ .

**Theorem 2** For  $\lambda < 0$ , for each closed cone  $(c_1, c_2, c_3)$  there is a **unique** AC shrinker defined for large  $t$  asymptotic to the given cone.

$\Rightarrow$  invariant AC shrinker ends are rigid.

## Constructing invariant non-steady AC ends

---

The problem of constructing an *invariant non-steady soliton AC end asymptotic to a given closed invariant cone* can be written as another singular initial value problem (SIVP) for the first-order ODE system: this time the SIVP is **irregular**.

**Theorem 1** For any  $\lambda \neq 0$ ,  $\exists$  a ! formal power series solution  $\mathcal{P}$  in  $t^{-1}$  determined by the cone;  $\exists$  a solution of the ODE system that is smooth in a nhd of  $t = +\infty$  and whose Taylor series is  $\mathcal{P}$ .

**Theorem 2** For  $\lambda < 0$ , for each closed cone  $(c_1, c_2, c_3)$  there is a **unique** AC shrinker defined for large  $t$  asymptotic to the given cone.

$\Rightarrow$  invariant AC shrinker ends are rigid.

**Theorem 3** Given  $\lambda > 0$  and any closed cone  $(c_1, c_2, c_3)$

- $\exists$  a 2-parameter family of AC soliton ends asymptotic to the given cone.
- Difference between two solutions is of order  $\exp(-\frac{\lambda}{6}t^2) * \text{polynomial}$ .
- If  $c_2 = c_3$ , then a 1-parameter subfamily has  $f_2 = f_3$ .

Flow lines of this  $4=(2+2)$ -parameter family of solutions fill an open subset of 5-dimensional phase space.  $\Rightarrow$  invariant AC expander ends are stable.

## Closed invariant $G_2$ cones

---

Helpful to analyse invariant  $G_2$ -structures on  $\mathbb{R}_+ \times \Sigma$  in terms of scale and homothety class of invariant metrics on  $\Sigma$ :

$$\text{scale } g := \sqrt[3]{f_1 f_2 f_3} = \sqrt[6]{\text{vol}(\Sigma)}$$

$$\text{homothety class } \frac{f_1}{g}, \frac{f_2}{g}, \frac{f_3}{g}$$

$\varphi$  closed and homothety class constant implies  $g$  linear and  $\varphi$  conical:

$$d\varphi = 0 \Rightarrow \frac{dg}{dt} = \frac{1}{6} \left( \frac{f_1^2}{g^2} + \frac{f_2^2}{g^2} + \frac{f_3^2}{g^2} \right) \Rightarrow f_i = c_i t$$

with

$$6c_1 c_2 c_3 = c_1^2 + c_2^2 + c_3^2. \quad (*)$$



## Closed invariant $G_2$ cones

---

Helpful to analyse invariant  $G_2$ -structures on  $\mathbb{R}_+ \times \Sigma$  in terms of scale and homothety class of invariant metrics on  $\Sigma$ :

$$\text{scale } g := \sqrt[3]{f_1 f_2 f_3} = \sqrt[6]{\text{vol}(\Sigma)}$$

$$\text{homothety class } \frac{f_1}{g}, \frac{f_2}{g}, \frac{f_3}{g}$$

$\varphi$  closed and homothety class constant implies  $g$  linear and  $\varphi$  conical:

$$d\varphi = 0 \Rightarrow \frac{dg}{dt} = \frac{1}{6} \left( \frac{f_1^2}{g^2} + \frac{f_2^2}{g^2} + \frac{f_3^2}{g^2} \right) \Rightarrow f_i = c_i t$$

with

$$6c_1 c_2 c_3 = c_1^2 + c_2^2 + c_3^2. \quad (*)$$

Note: *any* positive triple  $(c_1, c_2, c_3)$  can be uniquely rescaled to satisfy  $(*)$   
 $\rightsquigarrow$  2-parameter family of closed conical  $G_2$ -structures on  $\mathbb{R}_+ \times SU(3)/T^2$ .

In other words: **given a homothety class on  $\Sigma$ , there is a unique choice of “cone angle” that makes it a closed cone.**

## Understanding completeness

---

**Q:** Which of these invariant solitons extends to a complete AC solution?

## Understanding completeness

---

**Q:** Which of these invariant solitons extends to a complete AC solution?

*Tendency:*

If  $\frac{f_1}{g}$ ,  $\frac{f_2}{g}$ ,  $\frac{f_3}{g}$  remain bounded as  $t \rightarrow \infty$  then asymptotic to closed cone.

## Understanding completeness

---

**Q:** Which of these invariant solitons extends to a complete AC solution?

*Tendency:*

If  $\frac{f_1}{g}, \frac{f_2}{g}, \frac{f_3}{g}$  remain bounded as  $t \rightarrow \infty$  then asymptotic to closed cone.

**Rough strategy** for finding AC solitons on  $\Lambda_-^2 M = M \sqcup \mathbb{R}_+ \times \Sigma$ .

1. Solutions on  $(0, \epsilon) \times \Sigma$  that extend smoothly across  $M$  at  $t = 0$ ?
2. Solutions for large  $t$  asymptotic to prescribed closed cone  $(c_1, c_2, c_3)$ ?
3. Do solutions from **1** and **2** fit together?

## Comparison with other flows: the steady case

---

- All known *steady* solitons in Ricci flow have *sub-Euclidean* volume growth:
  - the Bryant soliton; Appleton's resolutions of (some of) its quotients.
  - Bryant soliton known to appear in a finite-time singularity of RF.
  - known Kähler examples have at most half-dimensional volume growth (Cao, Conlon–Deruelle). *Not* seen in finite-time singular behaviour of KRF.

## Comparison with other flows: the steady case

---

- All known *steady* solitons in Ricci flow have *sub-Euclidean* volume growth:
  - the Bryant soliton; Appleton's resolutions of (some of) its quotients.
  - Bryant soliton known to appear in a finite-time singularity of RF.
  - known Kähler examples have at most half-dimensional volume growth (Cao, Conlon–Deruelle). *Not* seen in finite-time singular behaviour of KRF.
- Our steady AC  $G_2$  solitons most closely resemble Joyce-Lee-Tsui's (JLT) *translating solitons* in Lagrangian mean curvature flow (LMCF).
  - Joyce conjectures JLT translating solitons *can* appear in finite-time singularities of LMCF if Floer homology is obstructed.
  - Speculate that our steady  $G_2$  solitons can also arise as finite-time singularities of Laplacian flow on a compact 7-manifold.

(Our 2-parameter family of AC  $G_2$  expanders on  $\Lambda^2 \mathbb{C}P^2$  resembles JLT's family of exact Maslov-zero LMCF expanders asymptotic to pairs of transverse Lagrangian 3-planes).

## Comparison with other flows: shrinkers

---

**Ricci flow:** One obvious significant difference: absence of *compact* shrinkers in  $G_2$  flow; associated with positive curvature in RF, whereas scalar curvature is non-positive for closed  $G_2$ -structures.

General theory for *noncompact complete shrinkers* in RF is well-developed:

- their properties are a hybrid of those of positively curved Einstein manifolds and spaces with non-negative Ricci, e.g. at most Euclidean volume growth.
- AC (gradient) shrinkers are extremely rigid—manifestation of parabolic backwards uniqueness phenomenon, also seen in MCF.
- AC end behaviour of our (highly symmetric)  $G_2$  shrinkers some indication such strong rigidity also holds for AC  $G_2$  (gradient?) shrinkers.

## Comparison with other flows: shrinkers

---

**Ricci flow:** One obvious significant difference: absence of *compact* shrinkers in  $G_2$  flow; associated with positive curvature in RF, whereas scalar curvature is non-positive for closed  $G_2$ -structures.

General theory for *noncompact complete shrinkers* in RF is well-developed:

- their properties are a hybrid of those of positively curved Einstein manifolds and spaces with non-negative Ricci, e.g. at most Euclidean volume growth.
- AC (gradient) shrinkers are extremely rigid—manifestation of parabolic backwards uniqueness phenomenon, also seen in MCF.
- AC end behaviour of our (highly symmetric)  $G_2$  shrinkers some indication such strong rigidity also holds for AC  $G_2$  (gradient?) shrinkers.

**LMCF:** self-shrinkers exist and do occur but *not* in the Maslov-zero (graded) setting. **Q:** Is there any natural condition to impose in the  $G_2$  setting that would rule out our AC shrinkers on  $\Lambda_+^2 \mathbb{S}^4$  and  $\Lambda_+^2 \mathbb{C}P^2$ ?

**KRF:** Feldman-Ilmanen-Knopf (FIK) constructed symmetric ALE Kähler shrinkers; simplest FIK shrinker does appear as a finite-time blowup of KRF on 1-point blowup of  $\mathbb{C}P^2$  and is associated with blowing down the point.



# Expanders

---

## Theorem (C)

For  $\lambda > 0$ , each  $\varphi_m$  extends to a complete solution with  $f_2 = f_3$ , and

$$\frac{f_i}{t} \rightarrow c_i$$

for  $(c_1, c_2, c_2)$  a closed cone with  $c_1 \leq c_2$ . These expanders are all smoothly asymptotic with rate  $-1$  to a unique invariant closed  $G_2$ -cone.

This gives 1-parameter families of expanders on both  $\Lambda_-^2 \mathbb{S}^4$  and  $\Lambda_-^2 \mathbb{C}P^2$ .

# Expanders

---

## Theorem (C)

For  $\lambda > 0$ , each  $\varphi_m$  extends to a complete solution with  $f_2 = f_3$ , and

$$\frac{f_i}{t} \rightarrow c_i$$

for  $(c_1, c_2, c_2)$  a closed cone with  $c_1 \leq c_2$ . These expanders are all smoothly asymptotic with rate  $-1$  to a unique invariant closed  $G_2$ -cone.

This gives 1-parameter families of expanders on both  $\Lambda_-^2 \mathbb{S}^4$  and  $\Lambda_-^2 \mathbb{C}P^2$ .

### *Very Strong Expectation*

- 1-1 correspondence with closed cones such that  $c_1 < c_2$ :  
i.e. any closed cone with  $c_2 = c_3$  on “one side” of the torsion-free cone  $(\frac{1}{2}, \frac{1}{2}, \frac{1}{2})$  is the AC end of a unique expander

# Expanders

---

## Theorem (C)

For  $\lambda > 0$ , each  $\varphi_m$  extends to a complete solution with  $f_2 = f_3$ , and

$$\frac{f_i}{t} \rightarrow c_i$$

for  $(c_1, c_2, c_2)$  a closed cone with  $c_1 \leq c_2$ . These expanders are all smoothly asymptotic with rate  $-1$  to a unique invariant closed  $G_2$ -cone.

This gives 1-parameter families of expanders on both  $\Lambda_-^2 \mathbb{S}^4$  and  $\Lambda_-^2 \mathbb{C}P^2$ .

### Very Strong Expectation

- 1-1 correspondence with closed cones such that  $c_1 < c_2$ :  
i.e. any closed cone with  $c_2 = c_3$  on “one side” of the torsion-free cone  $(\frac{1}{2}, \frac{1}{2}, \frac{1}{2})$  is the AC end of a unique expander

## Conjecture

For  $\lambda > 0$ , an **open** subfamily of  $\varphi_{m,c}$  (but not all) extend to complete solutions, defining a 2-parameter family of AC solitons on  $\Lambda_-^2 \mathbb{C}P^2$ .