

COHERENT at the Spallation Neutron Source

J. Adam,¹ P.S. Barbeau,^{1,2} P. Barton,³ A. Bolozdynya,⁴ B. Cabrera-Palmer,⁵ J.I. Collar,⁶ R. Cooper,³ R. Cooper,⁷ D. Dean,⁸ Y. Efremenko,^{4,9} S. Elliott,¹⁰ N. Fields,¹¹ M. Foxe,¹² A. Galindo-Uribarri,^{8,9} M. Gerling,⁵ M. Green,⁸ G. Greene,^{8,9} D. Hornback,⁸ T. Hossbach,¹² E.B. Iverson,⁸ S. Klein,³ A. Khromov,⁴ A. Kumpan,⁴ W. Lu,⁸ D. Markoff,^{13,2} P. Mueller,⁸ M. McIntyre,¹⁴ J. Newby,⁸ J. Orrell,¹² S. Penttila,⁸ G. Perumpilly,⁶ D. Radford,⁸ H. Ray,¹⁴ J. Raybern,^{1,8} D. Reyna,⁵ G. Rich,² D. Rimal,¹⁴ K. Scholberg*,^{1,†} B. Scholz,⁶ S. Suchyta,¹⁵ R. Tayloe,⁷ K. Vetter,^{15,3} and C.H. Yu⁸

¹*Department of Physics, Duke University, Durham, NC 27708, USA*

²*Triangle Universities Nuclear Laboratory, Durham, North Carolina, 27708, USA*

³*Lawrence Berkeley National Laboratory, Berkeley, CA 94720, USA*

⁴*National Research Nuclear University MEPhI, Moscow, 115409, Russia*

⁵*Sandia National Laboratories, Livermore, CA 94550, USA*

⁶*Enrico Fermi Institute, Kavli Institute for Cosmological Physics and
Department of Physics, University of Chicago, Chicago, IL 60637*

⁷*Department of Physics, Indiana University, Bloomington, IN, 47405, USA*

⁸*Oak Ridge National Laboratory, Oak Ridge, TN 37831, USA*

⁹*Department of Physics and Astronomy, University of Tennessee, Knoxville, TN 37996, USA*

¹⁰*Los Alamos National Laboratory, Los Alamos, NM, USA, 87545, USA*

¹¹*Enrico Fermi Institute, Kavli Institute for Cosmological Physics and
Department of Physics, University of Chicago, Chicago, IL 60637, USA*

¹²*Pacific Northwest Laboratory, Richland, WA 99352, USA*

¹³*Physics Department, North Carolina Central University, Durham, North Carolina 27707, USA*

¹⁴*Department of Physics, University of Florida, Gainesville, FL 32611, USA*

¹⁵*Department of Nuclear Engineering, University of California, Berkeley, CA, USA*

The measurement of coherent elastic neutrino-nucleus scattering can address a number of scientific questions called out in the 2007 Nuclear Physics Long Range Plan related to the nature of neutrinos and their role in supernova explosions. Ultimately, it can address questions of nuclear structure. In this position paper, we describe how the Spallation Neutron Source (SNS) at Oak Ridge National Laboratory, Tennessee, can be used for a measurement of coherent elastic neutrino-nucleus scattering, and describe the physics reach of different phases of such an experimental program, called COHERENT. The SNS provides an intense flux of neutrinos in the few tens-of-MeV range, with a sharply-pulsed timing structure that is beneficial for background rejection and highly advantageous for measurement of coherent elastic neutrino-nucleus scattering.

I. EXECUTIVE SUMMARY

The Spallation Neutron Source (SNS) offers unique opportunities to address 2007 Nuclear Physics Long Range Plan [1] questions, such as “What is the nature of the neutrinos, what are their masses, and how have they shaped the evolution of the universe?,” as well as questions related to nuclear reactions occurring in supernova explosions and the origin of the elements. Properties of the SNS as a neutrino source and a wide variety of opportunities are described in [2, 3].

A very interesting possibility for a stopped-pion neutrino source like the SNS is the detection of nuclear recoils from coherent elastic neutrino-nucleus scattering, which is within the reach of the current generation of low-threshold detectors [4]. This reaction has important implications for a range of nuclear physics problems including supernova processes and detection. Coherent elastic neutrino-nucleus scattering provides a new channel for measurement of the weak mixing angle, and also has excellent prospects for standard model tests; even a first-generation experiment has sensitivity beyond the current best limits on non-standard interactions of neutrinos and quarks. Sterile neutrino oscillation searches are also possible [5]. Beyond neutrino physics, a coherent elastic neutrino-nucleus scattering may eventually provide a new probe of nuclear structure [6, 7].

* Corresponding author

†Electronic address: schol1@phy.duke.edu

A stopped-pion beam has several advantages; the relatively high energies enhance the cross section, and recoil energies in the few to few tens of keV range bring detection within reach of the current generation of low-threshold detectors, which are scalable to relatively large target masses. Furthermore, the pulsed nature of the source allows both background reduction and precise characterization of the remaining background by measurement during the beam-off period. Finally, the different flavor content of the SNS flux (ν_μ , ν_e and $\bar{\nu}_\mu$) means that physics sensitivity is complementary to that for reactors ($\bar{\nu}_e$ only).

We envision three experimental phases with different experimental scales that will address different physics:

Phase 1: a few to few tens of kg of target material (depending on distance to the source) could make the first measurement.

Phase 2: a few tens to hundreds of kg of target material could set significant limits on non-standard neutrino interactions, and could also begin to address sterile neutrino oscillations, depending on configuration.

Phase 3: a tonne-scale or more experiment could begin to probe neutron distributions.

We have developed plans for a first phase measurement at the SNS, in the form of a program called COHERENT [8], for which we have identified three short-term detector possibilities, exploiting the low-threshold, low-background expertise of the NP community: CsI[Na] [9], germanium point-contact technology (e.g., [10]), and two-phase liquid xenon [11]. We aim to deploy as many types of detectors as resources will permit, and others are possible in the longer term. The rapid testing of different techniques by COHERENT will inform a future experimental plan. Following first results of a background measurement campaign (which is ongoing), a neutron-quiet location in the SNS basement at about a 20 meter distance from the neutrino source has been identified. A short-term goal is to measure neutrino-induced neutrons on lead, for which only theoretical calculations exist [12, 13], which may be a non-negligible background for detector setups employing lead as a shielding material and which is of intrinsic interest for, e.g., supernova detection [14]. Future phases of COHERENT may consider a pit located outside of the SNS target building to further reduce beam-associated backgrounds, as well as providing improved shielding from cosmogenic backgrounds.

II. INTRODUCTION

Coherent neutral-current (NC) neutrino-nucleus scattering (CENNS) was first predicted theoretically in 1974 [15] but has never been observed experimentally. The condition of coherence requires sufficiently small momentum transfer to the nucleon so that the waves of off-scattered nucleons in the nucleus are all in phase and contribute coherently. While interactions for neutrino energies in MeV to GeV range have coherent properties, neutrinos with energies less than 50 MeV are most favorable, as they largely fulfill the coherence condition in most target materials with nucleus recoil energy of tens of keV. The elastic NC interaction in particular leaves no observable signature other than low-energy nuclear recoils. Technical difficulties involved in the development of a large-scale, low-energy-threshold and low-background detectors have hampered the experimental realization of the CENNS measurement for more than three decades. However, recent innovations in dark-matter detector technology (e.g., [16, 17]) have made the unseen CENNS reaction testable. A well-defined neutrino source is the essential component for measurement of CENNS. The energy range of the stopped-pion Spallation Neutron Source (SNS) neutrinos is below 50 MeV, which is the optimal energy to observe pure coherent νA scattering. The detection is within the reach of the current generation of low-threshold detectors [4]. In addition to being a precise standard-model test, this reaction is also important for supernova processes and detection. Physics reach is described in more detail in Section IV and reference [4].

This document describes a program to exploit the SNS neutrinos for CENNS physics, to be implemented by the newly-formed COHERENT collaboration. The first phase would aim for a first measurement. Subsequent phases (possibly sharing resources with other initiatives for neutrino physics at the SNS [18, 19]) would aim for new tests of the standard model (SM); later phases could probe questions in nuclear physics. This document draws from a more comprehensive white paper covering opportunities with neutrinos at the SNS [2], and discusses recent progress and prospects.

III. THE SNS AS A NEUTRINO SOURCE

The SNS is the world's premier facility for neutron-scattering research, delivering approximately 6000 MW-hours of 1-GeV protons to the liquid mercury target in 700-ns-wide pulses at 60 Hz (more than 10^{14} protons/20 kJ per pulse).

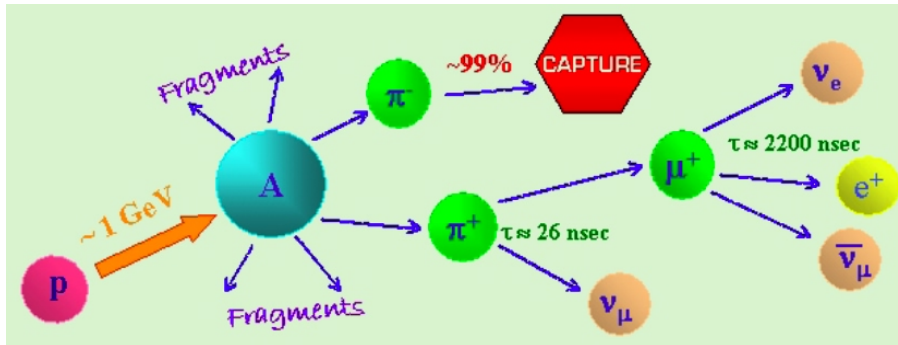


FIG. 1: SNS neutrino production mechanism.

Neutrons produced in spallation reactions with the mercury thermalize in hydrogenous moderators surrounding the target and are delivered to neutron-scattering instruments in the SNS experiment hall.

As a by-product, the SNS also provides the world's most intense pulsed source of neutrinos in the energy regime of interest for particle and nuclear astrophysics. Interactions of the proton beam in the mercury target produce mesons in addition to neutrons. These stop inside the dense mercury target and their subsequent decay chain, illustrated in Fig. 1, produces neutrinos with a flux of $\sim 2 \times 10^7 \text{ cm}^{-2}\text{s}^{-1}$ for all flavors at 20 m from the spallation target. This exceeds the neutrino flux at ISIS (where the KARMEN experiment was located [20]) by more than an order of magnitude.

The energy spectra of SNS neutrinos are shown in the right-hand panel of Fig. 2. These spectra are well known because almost all neutrinos come from decay-at-rest processes in which the kinematics are well defined. The decay of stopped pions produces monoenergetic muon neutrinos at 30 MeV. The subsequent three-body muon decay produces a spectrum of electron neutrinos and muon antineutrinos with energies up to 52.6 MeV.

The time structure of the SNS beam is particularly advantageous for neutrino studies. Time correlations between candidate events and the SNS proton beam pulse will greatly reduce background rates. As shown in the top left panel of Fig. 2, nearly all neutrinos will arrive within several microseconds of the 60-Hz proton beam pulses. As a result, background events resulting from cosmic rays will be suppressed by a factor of ~ 2000 by ignoring events which occur too long after a beam pulse. At the beginning of the beam spill, the neutrino flux is dominated by muon neutrinos resulting from pion decay, in principle making it possible to isolate pure NC events, since the ν_μ in the source have energies below the charged-current (CC) threshold. Backgrounds from beam-induced high-energy neutrons are present, but can be mitigated by appropriate siting and shielding. Beam-induced neutron backgrounds are greatly suppressed for $t \gtrsim 1 \mu\text{s}$ after the start of the beam spill, while the neutrino production, governed by the muon lifetime ($\tau_\mu \sim 2.2 \mu\text{s}$), proceeds for several microseconds. This time structure presents a great advantage over a long-duty-factor facility such as the Los Alamos Neutron Science Center (LANSCE), where the LSND experiment was located [21]. Figure 3 shows the expected fluence at the SNS, compared to what would be expected for a nearby supernova.

In general, one wants high neutrino flux (with flux roughly proportional to proton beam power), sharp pulses to enable rejection of off-beam background, and well-understood neutrino spectra. Ideally pulses should be shorter than the muon-decay lifetime, and separated by at least several τ_μ . Proton energies are ideally less than about 1 GeV in order to minimize contamination from, e.g., kaons and decay-in-flight pions, which can be produced at a significant rate at higher proton beam energies. Proton energies and target configuration resulting in a high fraction of pion decays at rest will lead to a clean decay-at-rest spectrum and well-known flavor composition. The SNS satisfies all of these requirements, and overall is the only facility that within the next decade can provide 1-MW-level intensity, short duty factor, clean decay-at-rest neutrinos. We note that a second SNS target station may eventually be built; this would provide additional flux, although specific timing characteristics are as yet unknown.

IV. PHYSICS MOTIVATIONS

A. Coherent Elastic Neutrino-Nucleus Scattering

The cross section for a spin-zero nucleus, neglecting radiative corrections, is given by [24],

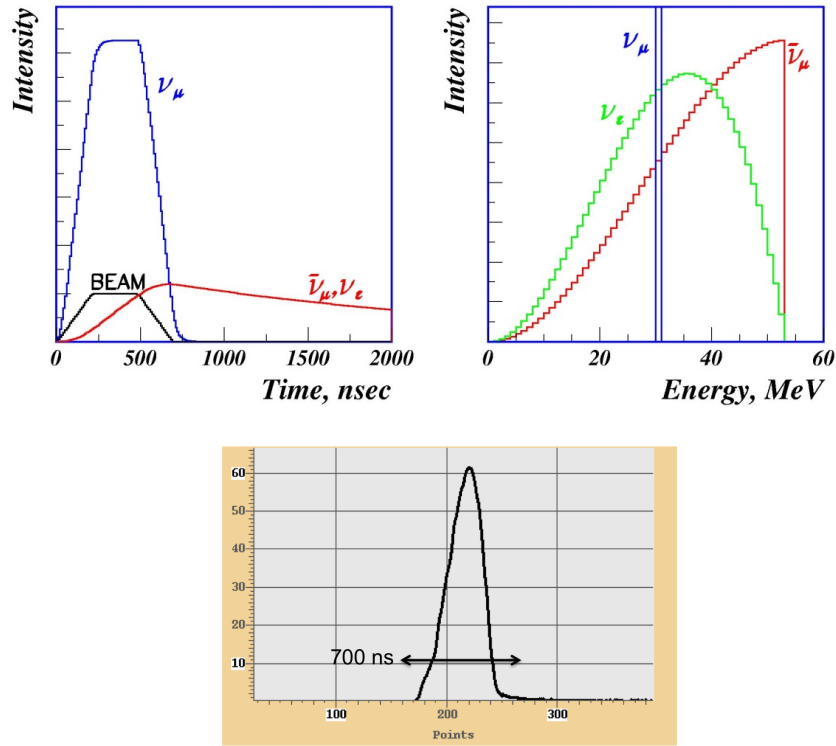


FIG. 2: Time and energy distributions for the different neutrino flavors produced at the SNS. The top plots are from a 2003 study [22]; the bottom plot shows proton beam structure on target from recent SNS running [23]. The prompt (pion-decay) component of the neutrino flux should closely follow the proton time structure.

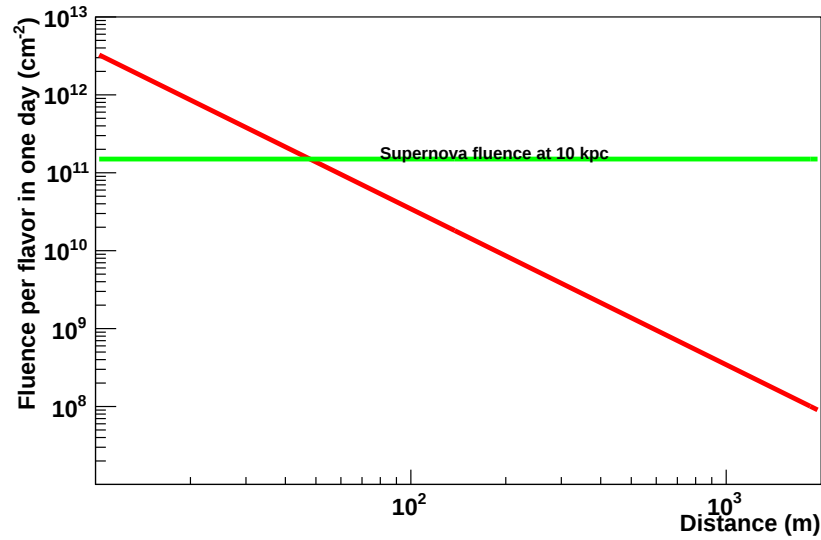


FIG. 3: The red line shows integrated fluence per flavor in one day for the SNS neutrino flux as a function of distance from the source. The green solid line shows approximate fluence per flavor for a supernova at 10 kpc for the full burst.

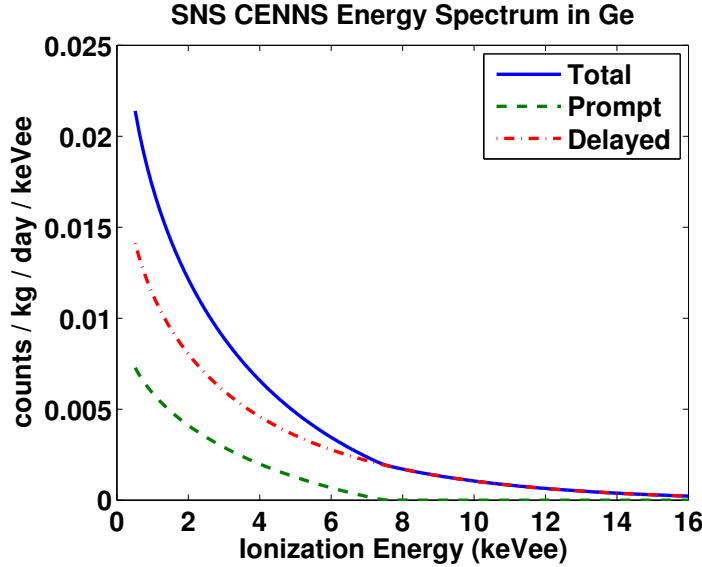


FIG. 4: CENNS ionization-energy (keVee) spectrum (events per keVee per kg active mass of Ge per day of observation time) in a Ge detector located 20 m from the SNS target. The prompt and delayed components of the spectrum are shown separately.

$$\frac{d\sigma}{dT}(E, T) = \frac{G_F^2}{2\pi} M \left[2 - \frac{2T}{E} + \left(\frac{T}{E} \right)^2 - \frac{MT}{E^2} \right] \frac{Q_W^2}{4} F^2(Q^2). \quad (1)$$

where E is the incident neutrino energy, T is the nuclear recoil energy, M is the nuclear mass, F is the ground-state elastic form factor, Q_w is the weak nuclear charge, and G_F is the Fermi constant. The condition for coherence requires that momentum transfer $Q \lesssim \frac{1}{R}$, where R is the nuclear radius. This condition is largely satisfied for neutrino energies up to ~ 50 MeV for medium A nuclei. Typical values of the total coherent elastic cross section are in the range $\sim 10^{-39}$ cm², which is relatively high compared to other neutrino interactions in this energy range (e.g., CC inverse beta decay on protons has a cross section $\sigma_{\bar{\nu}_e p} \sim 10^{-40}$ cm², and elastic neutrino-electron scattering has a cross section $\sigma_{\nu_e e} \sim 10^{-43}$ cm²).

Although ongoing efforts to observe CENNS at reactors [25–27] are promising, a stopped-pion beam has several advantages with respect to the reactor experiments. Higher recoil energies bring detection within reach of the current generation of low-threshold detectors which are scalable to relatively large target masses. Furthermore, the pulsed nature of the source (see Fig. 2) allows both background reduction and precise characterization of the remaining background by measurement during the beam-off period. Finally, the different flavor content ($\nu_e, \nu_\mu, \bar{\nu}_\mu$) of the SNS flux means that physics sensitivity is complementary to that for reactors, which provide only $\bar{\nu}_e$; for example, non-standard neutrino interactions (NSI) may be flavor-dependent. The expected rates for the SNS are quite promising for relevant low-energy detector materials [4]: see Fig. 4.

CENNS reactions are important in stellar core-collapse processes [15], as well as being useful for core-collapse supernova neutrino detection [24]. A rate measurement will have bearing on supernova neutrino physics. The CENNS cross section is predicted by the SM, and form-factor uncertainties are small [24]. Therefore a measured deviation from prediction could be a signature of new physics. We also note that successful measurement of CENNS in the energy range of solar and atmospheric neutrinos will be immediately useful for direct dark matter search experiments, for which solar and atmospheric neutrinos will eventually represent a background [28, 29].

We note also that the measured neutrino flux will be a valuable input to the proposed OscSNS experiment [19], as well as the CAPTAIN [18] experiment if sited at the SNS, and CENNS could potentially be used for sterile neutrino oscillation searches (e.g., [5]).

A few of the possible physics measurements are described in more detail below.

B. Standard Model Tests With Coherent Scattering

According to Eq. 1, the SM predicts a coherent elastic scattering rate proportional to Q_w^2 , the weak charge given by $Q_w = N - (1 - 4\sin^2\theta_W)Z$, where Z is the number of protons, N is the number of neutrons, and θ_W is the weak mixing angle. Therefore the weak mixing angle can be extracted from the measured absolute cross section, at a typical Q value of 0.04 GeV/ c . A deviation from the SM prediction could indicate new physics. If the absolute cross section can be measured to 10%, there will be an uncertainty on $\sin^2\theta_W$ of $\sim 5\%$. One might improve this uncertainty by looking at ratios of rates in targets with different N and Z , to cancel common flux uncertainties; future use of enriched neon is a possibility. There are existing precision measurements from atomic parity violation [30, 31], SLAC E158 [32] and NuTeV [33]. However there is no previous neutrino scattering measurement in this region of Q . This Q value is relatively close to that of the Qweak parity-violating electron scattering experiment at JLAB [34, 35]. However CENNS tests the SM in a different channel and therefore is complementary.

In particular, such an experiment can search for non-standard interactions (NSI) of neutrinos with nuclei. Existing and planned precision measurements of the weak mixing angle at low Q do not constrain new physics that is specific to neutrino-nucleon interactions. Reference [36] explores the sensitivity of a CENNS experiment on the ton scale to some specific physics beyond the standard model, including models with extra neutral gauge bosons, leptoquarks and R-parity breaking interactions.

The signature of NSI is a deviation from the expected cross section. Reference [4] explores the sensitivity of an experiment at the SNS. As shown in the reference, under reasonable assumptions, if the rate predicted by the SM is observed, neutrino scattering limits more stringent than current ones [37, 38] by about an order of magnitude can be obtained for some parameters.

Searches for NSI are based on precise knowledge of the nuclear form factors, which are known to better than 5% [24], so that a deviation from the SM prediction larger than that would indicate physics beyond the SM.

Another possibility is to use CENNS to look for anomalous neutrino magnetic moment [4], for which the signature would be a distortion of the low-energy recoil spectrum [39]. The existing limits on magnetic moment for muon neutrinos are relatively weak [40], and the muon-flavor content of the stopped-pion neutrino flux would enable improvement of these limits. Such measurements require very good understanding of detector energy response at low energy.

C. Nuclear Structure from Coherent Scattering

If we assume that the SM is a good description of nature, then with sufficient precision one can measure neutron form factors. This physics could be within reach of a next-generation coherent scattering experiment. One of the basic properties of a nucleus is its size, or radius, typically defined as

$$\langle R_{n,p}^2 \rangle^{1/2} = \left(\frac{\int \rho_{n,p}(r) r^2 d^3r}{\int \rho_{n,p}(r) d^3r} \right)^{1/2}, \quad (2)$$

where $\rho_{n,p}(r)$ are the neutron and proton density distributions. Proton distributions in nuclei have been measured in the past to a high degree of precision. In contrast, neutron distributions are still poorly known. A measurement of neutron distributions could have an impact on a wide range of fields, from nuclear physics to astrophysics.

Previous measurements of the neutron radius have used hadronic scattering, and result in uncertainties of about $\sim 10\%$ [41]. A new measurement, being done at Jefferson Laboratory by the PREX experiment, uses parity-violating electron scattering to measure the neutron radius of lead. The current uncertainty in the neutron radius from this experiment is about 2.5% [42]. An alternate method, first suggested in [6], is to study the neutron radius through neutrino-nucleus coherent scattering. References [7, 43, 44] explore this possibility. This measurement is very challenging, as it requires very large statistics and excellent understanding of detector energy response; however it is conceivable for a next-generation CENNS experiment.

V. EXPERIMENTAL OPPORTUNITIES

A. Detectors to Measure Coherent Elastic νA Scattering

We envision approximately three experimental phases on different scales that will address different physics:

Phase 1: a few to few tens of kg of target material (depending on distance to the source) could make the first measurement.

Phase 2: a few tens to hundreds of kg of target material could set significant limits on NSI, and could also begin to address sterile neutrino oscillations, depending on configuration.

Phase 3: a ton-scale or more experiment could begin to probe neutron distributions and neutrino magnetic moment.

Various technologies are suitable at different scales. The COHERENT collaboration [45]¹ has identified three possibilities for prompt deployment: CsI[Na] [9], germanium point-contact technology (e.g., [10]), and two-phase xenon [11]. Conceivably more than one technology could be deployed simultaneously, depending on available resources, siting and background issues. We aim to deploy as many types as resources permit. There are also possibilities beyond those discussed here (e.g., single-phase argon [46, 47], bolometers [48]).

The following subsections describe these different possibilities.

1. CsI[Na] Detectors

CsI[Na] scintillators present several advantages for CENNS neutrino measurements. These are briefly listed below:

- The large CENNS cross-section from both recoiling species, cesium and iodine, provides ~ 800 recoils per 15 kg of CsI[Na] per year above the expected ~ 5 keVnr threshold of this detector. Both recoiling species are essentially indistinguishable due to their very similar mass, greatly simplifying understanding the response of the detector.
- The quenching factor for nuclear recoils in this material over the energy region of interest has been carefully characterized, using the methods described in [49] (Fig. 5). Its value is sufficiently large to expect a realistic ~ 5 keVnr threshold.
- It should be possible to perform statistical discrimination (as opposed to event-by-event discrimination) between nuclear and electron recoils at the level of $\sim 1,000$ accumulated events. This is based on a difference of ~ 60 ns between the scintillation decay times observed for these two families of events [50]. Similar differences have already been exploited to implement this discrimination in NaI[Tl] scintillators dedicated to WIMP detection.
- Prototype crystals grown from screened salts containing low levels of U, Th, ^{40}K , ^{87}Rb , and $^{134,137}\text{Cs}$ should provide a neutrino signal-to-background ratio of $\mathcal{O}(5)$, even prior to (anticoincident) background subtraction.
- Several other practical advantages exist: CsI[Na] exhibits a high light yield (64 photons/keVee) and has the best match to the response curve of alkali photomultipliers of any scintillator. It is a rugged room-temperature detector material, and is also relatively inexpensive ($\sim \$1/\text{g}$), permitting an eventual increase in mass to a ~ 100 kg target, which would allow to explore the most interesting physics goals planned for COHERENT. CsI[Na] lacks the excessive afterglow (phosphorescence) that is characteristic of CsI[Tl] [50], an important feature in a search involving small scintillation signals in a detector operated at ground level.

A 15 kg CsI[Na] detector (CoSI) and its shielding (Fig. 1), built by University of Chicago and Pacific Northwest National Laboratory COHERENT collaborators, is being completed at the time of this writing, aiming at installation at the SNS during 2014. This device should be sufficient for a first measurement of CENNS at the SNS, while providing operational experience towards a final array of crystals with total mass $\mathcal{O}(100)$ kg.

2. Germanium Detectors

A powerful approach to detecting CENNS at the SNS is to use point contact High-Purity Germanium (HPGe) detectors. Point contact HPGe detectors feature very small readout electrodes that offer extremely low electronic noise [51]. This low electronic noise in turn offers excellent energy resolution and detection thresholds below 1 keV. P-type point contact (PPC) detectors with an electronic noise on the order of 150 eV FWHM have been demonstrated [16, 52]. The excellent energy resolution afforded by PPC detectors also results in low systematic uncertainties associated with background subtraction.

¹ Note the formerly-proposed collaboration name “CSI (Coherent Scattering Investigations)” has now changed to COHERENT.

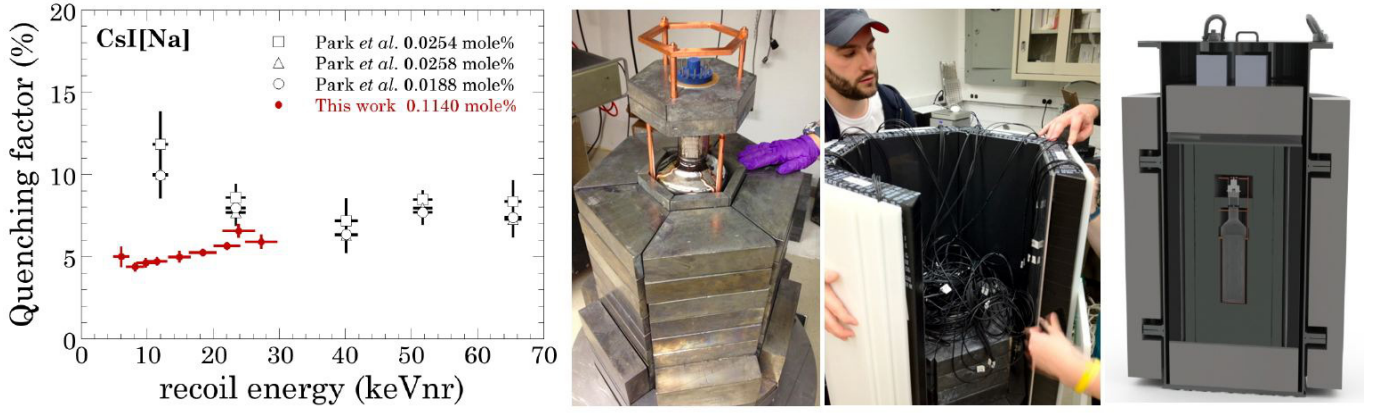


FIG. 5: *Left*: Quenching factor for low-energy nuclear recoils in CsI[Na], measured using the method described in [49]. The range of recoil energies is relevant for SNS neutrino energies. *Center panels*: Assembly of a 2 kg CsI[Na] prototype within the shield intended for SNS installation. The detector features selected salts screened against radioactive contaminants, OFHC copper parts, a low-background ET 9390UFL photomultiplier, archeological-quality lead shielding (<0.02 Bq/kg ^{210}Pb), a 99.6% efficient muon veto, and neutron moderator and absorber. *Right*: Full model of the shield containing the 15 kg detector, showing the steel container which could be used to lower the assembly into a shallow pit.

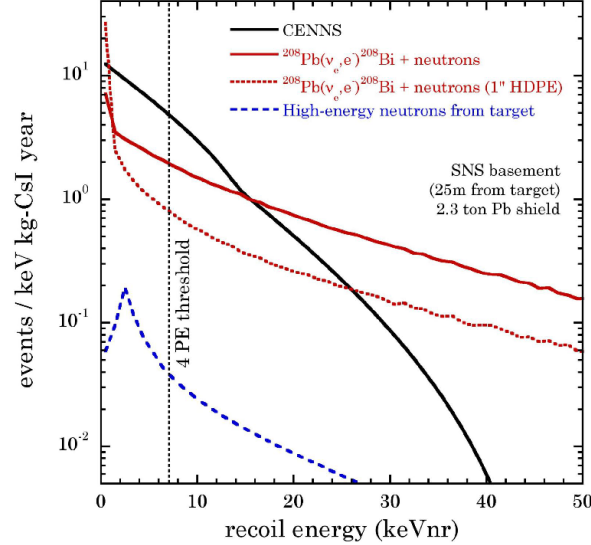


FIG. 6: From [9]: beam-associated neutrino signals and backgrounds expected in the 14-kg CsI[Na] detector at the SNS. Nuclear recoils induced by high-energy neutrons escaping the SNS target monolith can be sufficiently suppressed in a basement location discussed in the main text (the simulation shown here includes an electron recoil component from gamma emission following neutron thermalization). The ratio of CENNS events above threshold to those induced by the neutron emission following $^{208}\text{Pb}(\nu_e, e)^{208}\text{Bi}$ is 1.2:1. This ratio becomes 3.3:1 when the innermost 1 inch of Pb in the shielding is substituted by high-density polyethylene (HDPE). The possibility exists for a convincing simultaneous measurement of both neutrino cross-sections (see text).

HPGe detectors represent a mature and widely-used technology, with PPC detectors commercially available from, e.g., CANBERRA and ORTEC. The high intrinsic radiopurity of HPGe material, and the availability of low-background cryostats, make low-background operation a reality. Existing detectors have demonstrated intrinsic background levels suitable for the CENNS measurement [16]. The possibility of leveraging DOE NP supported technology developed for the MAJORANA DEMONSTRATOR [10], as well as components and detectors fabricated or procured for Demonstrator-related R&D activities, offer a potential cost-effective path to a deployable system. Detector mounting hardware (Fig. 7) and modular cryostats (Fig. 8) designed by the Majorana collaboration are capable of deploying germanium detectors in modular arrays ranging from single crystals to arrays with a total mass of 20 kg or more.

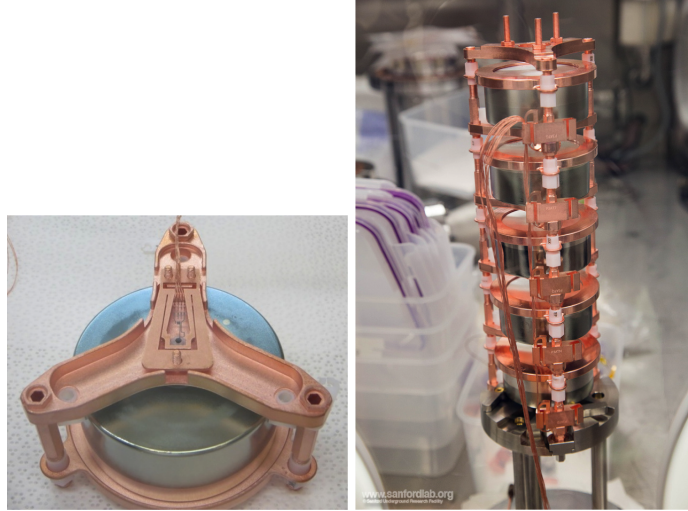


FIG. 7: Canberra Broad-Energy Germanium (BEGe) PPC-type detectors in the Majorana Demonstrator detector mounting hardware. Left: a single mounted detector with low-mass front end electronics. Right: A string of 5 detectors with total mass of ~ 3 kg (Photo credit: Matthew Kapust, Sanford Underground Research Facility).

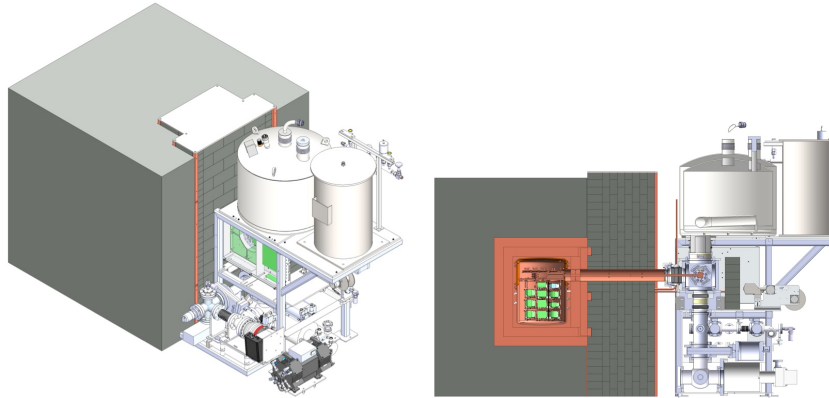


FIG. 8: A Majorana Demonstrator modular cryostat with capacity of ~ 20 kg of germanium detectors, deployed in a potential passive shield of copper and lead.

Figure 4 shows the spectrum of energy deposition, in terms of equivalent ionization energy (keVee) for a detector located 20 m from the SNS target, where the neutrino flux is $2 \times 10^7 \nu/\text{s}/\text{cm}^2$. The prompt and delayed components to the spectrum are shown separately. The cross-section form factor is modeled following [53] and the quenching factor in Ge following [54], which is well understood to be $20\% \pm 1\%$ near the low energy threshold of PPC detectors. At a distance of 20 m from the target, for a 1-keV ionization energy threshold, the signal is approximately 3.6 events/kg/month. A detector array with 20 kg of active mass, such as the one pictured in Fig. 8, would observe 400 signal events in less than 6 months.

By using a low-background detector and shielding, and taking advantage of the pulsed structure of the SNS beam, the intrinsic detector backgrounds would be far below the expected signal rate. Because germanium is relatively dense ($5.3 \text{ g}/\text{cm}^3$), the detector will be compact, and a fairly simple scintillator-based active veto could reduce the cosmic-ray muon rate to far below background. Cosmic-spallation neutrons and beam-related neutrons are left as the challenging sources of background; efforts are underway to measure these neutron backgrounds and determine shielded germanium PPC detector response.

In the longer term, the excellent germanium energy resolution should also lend itself to more precise studies of CENNS, by adding more detectors to the array or deploying additional arrays.

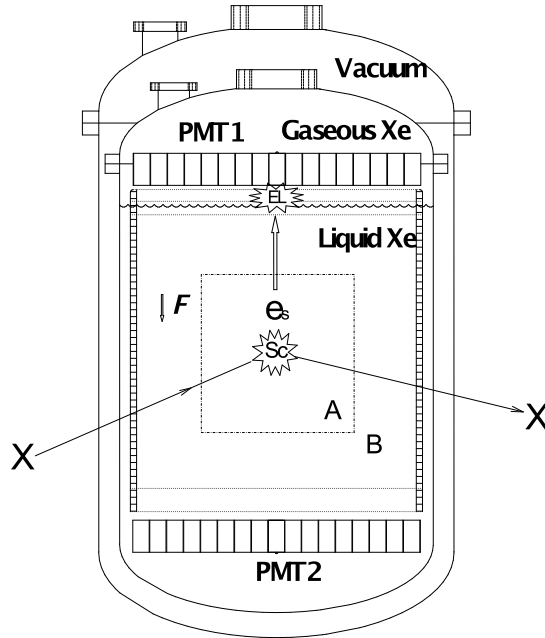


FIG. 9: Principle of operation for a wall-less liquid xenon emission detector detecting a hypothetical weakly-interacting particle X . Sc: scintillation flash generated at the point of primary interaction between X and Xe atoms; EL: electroluminescence flash of gaseous Xe excited by electrons extracted from liquid Xe by electric field F and drifting through the gas at high electric fields (>1 kV/cm/bar); PMT1 and PMT2: arrays of photodetectors detecting Sc and EL signals; A: the fiducial volume where events considered to be useful occur; B: the shielding layer of LXe. The active volume of the detector is surrounded with highly reflective cylindrical PTFE reflector embodied with drift electrode structure providing a uniform field F . The detector is enclosed in a vacuum cryostat made of low-background pure titanium.

3. Two-Phase Xenon Detectors

Since 1992, the employment of a liquid xenon detector has been considered for neutrino magnetic moment searches by probing $\bar{\nu}_e$ -electron scattering cross sections for deposited energies below 100 keV. Such an experiment could be performed with a moderately-sized (~ 1 ton mass) LXe emission detector [55].

The emission method of particle detection invented 40 years ago at the Moscow Engineering Physics Institute (MEPhI) Department of Experimental Nuclear Physics [56] allows an arrangement of a “wall-less” detector sensitive to single ionization electrons [57]. See Fig. 9. A detector using this technology is also sensitive to CENNS-induced recoils.

To measure the ionization yield of heavy nuclear recoils, an experiment is underway that will model the detection of xenon nuclear recoils through a study of the elastic scattering of a monochromatic filtered beam of neutrons from the IRT MEPhI reactor with a 5-kg LXe emission detector recently used for detection of single electrons [11, 58–60].

For observation of CENNS at the SNS, we consider the emission detector RED100 with a 100-kg LXe working medium. The basic parameters for construction of this detector are the following: a low-background titanium cryostat, a readout system based on two arrays of low-background PMTs (e.g., Hamamatsu R11410) located in the gas phase and in the liquid below the grid cathode, and a PTFE light-collection system embodied with a drift electrode system as shown in Fig. 9, a cryogenic system based on thermo-siphon technology similar to that used in the LUX detector, and a location in a borehole 10 m underground at 40 m distance from the target. For these conditions, the total expected event rate is 1470 events/year.

B. Siting Possibilities

Several potential sites have been identified inside the target building at which a CENNS experiment could be deployed within the next few years. They are shown in Fig. 10. One of these (site 4) is in a basement with about

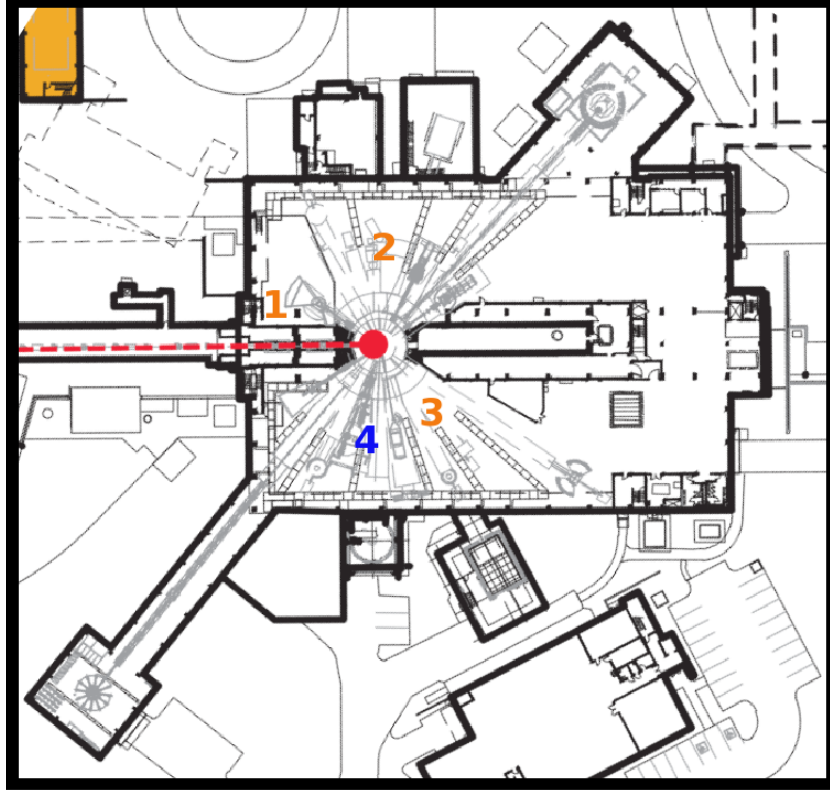


FIG. 10: Potential sites for deployment of a CENNS experiment at the SNS.

8 meters water-equivalent of overburden; the others are on the SNS target building floor level. At these locations, detectors can be deployed at distances between around 15-35 meters from the SNS source. Additional shielding can be deployed at the surface locations. As described in the next section, site 4 shows very low background levels and is our choice for the first set of measurements.

A site outside the SNS target building is also possible, and might be the best choice for a later phase, larger detector, although is less desirable due to farther distance from the neutrino source. A detector could potentially be deployed in a borehole for shielding. However it will be more expensive to deploy an experiment outside the SNS target building due to extra construction costs.

C. Background Measurements

Understanding and reduction of backgrounds are critical for the COHERENT experimental program. Steady-state backgrounds, such as radioactivity and cosmogenics, will be reduced by the SNS duty factor ($10^3 - 10^4$), and may or may not need additional mitigation with shielding and cleanliness of detector materials, depending on siting and detector properties. Steady-state backgrounds can also be very well understood using data taken outside the SNS beam window. However, beam-related backgrounds, especially fast neutrons, will have to be very carefully studied, using ancillary measurements and modeling, and shielded against. These beam-related backgrounds will likely be highly site-dependent, and also possibly time-dependent.

A background measurement campaign has been underway since fall of 2013 (in collaboration also with SNS personnel for whom neutron background measurements are independently useful). We are characterizing the backgrounds at the candidate locations (and beyond) and during different SNS running conditions. We have employed a two-plane neutron scatter camera and a liquid scintillator array deployed in different configurations, as well as a low-background point-contact germanium detector.

So far we have results from a measurements at site 2, between SNS beamlines 13 and 14 on the target hall 4, as well as measurements at different locations along a hallway parallel to the beamline in the basement around site 4. While there appears to be a substantial flux of neutrons in time with the beam pulse at site 2, such that a CENNS

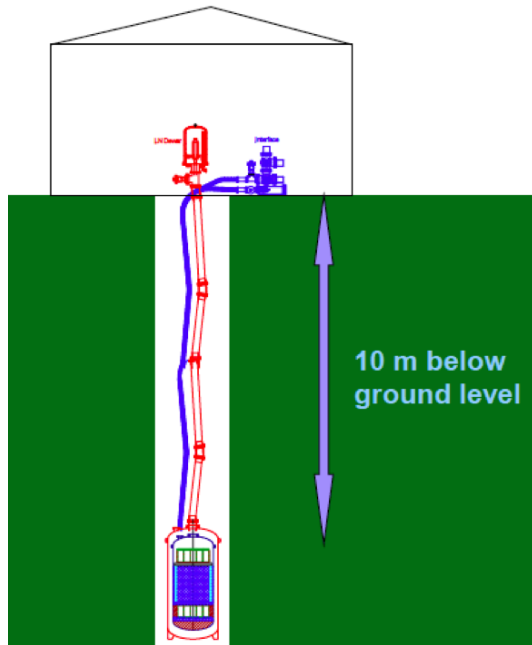


FIG. 11: Borehole concept for detector deployment.

measurement would be difficult, the basement is very neutron-quiet. Results are shown in Fig. 12. We therefore favor site 4, near positions 4 and 5, about 20 m from the SNS target, for the first measurements described in Section VI and possibly for the first CENNS measurement. However, we are continuing to explore different sites.

VI. NEUTRINO-INDUCED NEUTRONS

In the course of our studies, we have understood that neutrino-induced neutrons (NINs) on shielding material such as lead can be a non-negligible background for coherent scattering; see, e.g., Fig. 6 from reference [9]. The measurement of this cross section is of intrinsic interest, as is a supernova neutrino detection mechanism [12, 14], currently used for the HALO experiment [61]². Although some theoretical calculations exist [12, 13], these cross sections have never been measured. We are pursuing an opportunity to measure these cross sections on lead and potentially others (e.g., iron, copper) using the shielding setup from [9] holding liquid scintillator detectors, and additionally an easily-assembled design comprising scintillators surrounded by lead (or other material of interest) and water, shown in Fig. 13.

VII. CONCLUSION

We have outlined a program for taking advantage of the extremely high-quality stopped-pion neutrino source available at the Spallation Neutron Source available at Oak Ridge National Laboratory for CENNS measurements, addressing a number of scientific questions called out in the 2007 Nuclear Physics Long Range Plan.

² However solar-neutrino induced neutrons are not the source of the DAMA/LIBRA seasonal modulation proposed in [62]; see [63].

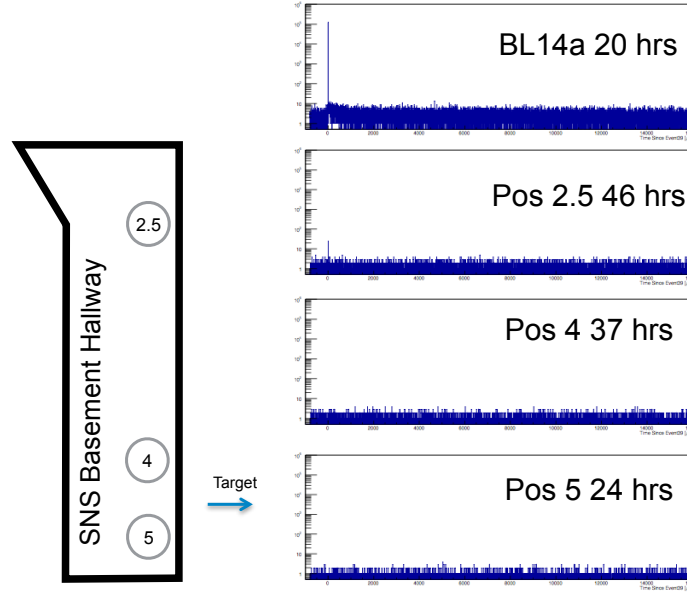


FIG. 12: Results of background measurements with a single-plane scintillator array in various locations. The four plots show the neutron signal rates as a function of time with respect to the beam pulse. The top plot is at the target-hall-floor site (site 2 on Fig. 10). The lower three plots are at the labeled positions in the basement (site 4 on Fig. 10). The SNS target is to the right of position 5. The neutron rates in time with the beam are very low at positions 4 and 5.

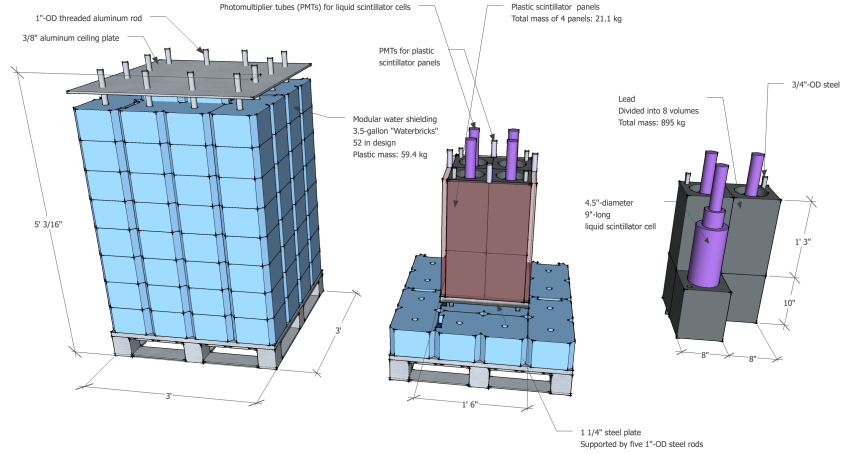


FIG. 13: Design for short-term measurement of NINs at the SNS.

Acknowledgments

We are grateful for excellent logistical support and advice from local SNS and ORNL personnel. Much of the background measurement work was done using ORNL SEED funds, as well as Sandia LDRD and NA-22 support.

-
- [1] (2008), 0809.3137.
 - [2] A. Bolozdynya *et al.*, (2012), 1211.5199.
 - [3] V. Cianciolo *et al.*, (2012), 1212.5190.
 - [4] K. Scholberg, Phys.Rev. **D73**, 033005 (2006), hep-ex/0511042.

- [5] A. Anderson *et al.*, Phys.Rev. **D86**, 013004 (2012), 1201.3805.
- [6] P. Amanik and G. McLaughlin, J.Phys.G **G36**, 015105 (2009).
- [7] K. Patton, J. Engel, G. McLaughlin, and N. Schunck, (2012), 1207.0693.
- [8] COHERENT Collaboration, D. Akimov *et al.*, (2013), 1310.0125.
- [9] J. Collar *et al.*, (2014), 1407.7524.
- [10] Majorana Collaboration, N. Abgrall *et al.*, Adv.High Energy Phys. **2014**, 365432 (2014), 1308.1633.
- [11] D. Y. Akimov *et al.*, JINST **8**, P05017 (2013), 1303.7338.
- [12] J. Engel, G. McLaughlin, and C. Volpe, Phys.Rev. **D67**, 013005 (2003), hep-ph/0209267.
- [13] E. Kolbe, K. Langanke, G. Martinez-Pinedo, and P. Vogel, J.Phys.G **G29**, 2569 (2003), nucl-th/0311022.
- [14] D. Vaananen and C. Volpe, (2011), 1105.6225.
- [15] D. Freedman, D. Schramm, and D. Tubbs, Ann.Rev.Nucl.Part.Sci. **27**, 167 (1977).
- [16] CoGeNT Collaboration, C. Aalseth *et al.*, Phys.Rev. **D88**, 012002 (2013), 1208.5737.
- [17] V. Chepel and H. Araujo, (2012), 1207.2292.
- [18] The CAPTAIN Collaboration, H. Berns *et al.*, (2013), 1309.1740.
- [19] M. Elnimr *et al.*, (2013), 1307.7097.
- [20] KARMEN, B. Zeitnitz, Prog.Part.Nucl.Phys. **32**, 351 (1994).
- [21] LSND, C. Athanassopoulos *et al.*, Nucl.Instrum.Meth. **A388**, 149 (1997), nucl-ex/9605002.
- [22] F. Avignone and Y. Efremenko, J.Phys. **G29**, 2615 (2003).
- [23] J. Galambos, SNS High Power Proton Beam Experience and Plans, Talk at COHERENT meeting.
- [24] C. Horowitz, K. Coakley, and D. McKinsey, Phys.Rev. **D68**, 023005 (2003), astro-ph/0302071.
- [25] P. Barbeau, J. Collar, and O. Tench, JCAP **0709**, 009 (2007), nucl-ex/0701012.
- [26] CoGeNT, J. Collar, J.Phys.Conf.Ser. **136**, 022009 (2008).
- [27] H. Wong, H.-B. Li, J. Li, Q. Yue, and Z.-Y. Zhou, J.Phys.Conf.Ser. **39**, 266 (2006), hep-ex/0511001.
- [28] J. Monroe and P. Fisher, Phys.Rev. **D76**, 033007 (2007), 0706.3019.
- [29] J. Billard, L. Strigari, and E. Figueroa-Feliciano, (2013), 1307.5458.
- [30] S. Bennett and C. Wieman, Phys.Rev.Lett. **82**, 2484 (1999), hep-ex/9903022.
- [31] Particle Data Group, S. Eidelman *et al.*, Phys.Lett. **B592**, 1 (2004).
- [32] SLAC E158, P. Anthony *et al.*, Phys.Rev.Lett. **95**, 081601 (2005), hep-ex/0504049.
- [33] NuTeV, G. Zeller *et al.*, Phys.Rev.Lett. **88**, 091802 (2002), hep-ex/01110059.
- [34] Qweak, W. van Oers, (2007), 0708.1972 [nucl-ex].
- [35] Qweak Collaboration, D. Androic *et al.*, (2013), 1307.5275.
- [36] J. Barranco, O. Miranda, and T. Rashba, Phys.Rev. **D76**, 073008 (2007), hep-ph/0702175.
- [37] CHARM, J. Dorenbosch *et al.*, Phys.Lett. **B180**, 303 (1986).
- [38] S. Davidson, C. Peña Garay, N. Rius, and A. Santamaria, JHEP **03**, 011 (2003), hep-ph/0302093.
- [39] P. Vogel and J. Engel, Phys.Rev. **D39**, 3378 (1989).
- [40] LSND, L. B. Auerbach *et al.*, Phys.Rev. **D63**, 112001 (2001), hep-ex/0101039.
- [41] C. Horowitz, S. Pollock, P. Souder, and R. Michaels, Phys.Rev. **C63**, 025501 (2001), nucl-th/9912038.
- [42] S. Abrahamyan *et al.*, Phys.Rev.Lett. **108**, 112502 (2012), 1201.2568.
- [43] P. Amanik and G. McLaughlin, (2007), 0707.4191 [hep-ph].
- [44] K. Patton, G. McLaughlin, and K. Scholberg, Int.J.Mod.Phys. **E22**, 1330013 (2013).
- [45] (2014), <http://coherent.phy.duke.edu>.
- [46] K. Scholberg *et al.*, (2009), 0910.1989.
- [47] S. Brice *et al.*, Phys.Rev. **D89**, 072004 (2014), 1311.5958.
- [48] J. A. Formaggio, E. Figueroa-Feliciano, and A. Anderson, Phys.Rev. **D85**, 013009 (2012), 1107.3512.
- [49] J. Collar, (2013), 1302.0796.
- [50] CosI collaboration, In preparation.
- [51] P. N. Luke, F. Goulding, N. Madden, and R. Pehl, IEEE Trans.Nucl.Sci. **36**, 926 (1989).
- [52] C. Aalseth *et al.*, Phys.Rev.Lett. **107**, 141301 (2011), 1106.0650.
- [53] S. Klein and J. Nystrand, Phys.Rev. **C60**, 014903 (1999), hep-ph/9902259.
- [54] P. Barbeau, (2009), Ph.D. thesis, AAT-3386971.
- [55] M. Baldo-Ceolin *et al.*, Large volume liquid Xe detector for anti-electron-neutrino magnetic moment measurement from anti-electron-neutrino e⁻ scattering at nuclear power reactor, in *Proceedings of Dubna 1992, Weak and electromagnetic interactions in nuclei*, pp. 753–759, 1992.
- [56] B. Dolgoshein, V. Lebedenko, and B. Rodionov, JETP Lett. **11**, 351 (1970).
- [57] A. Bolozdynya, V. Egorov, B. Rodionov, and V. Miroshnichenko, IEEE Trans.Nucl.Sci. **42**, 565 (1995).
- [58] http://enpl.mephi.ru/?is_eng=1.
- [59] V. Solovov *et al.*, IEEE Nucl.Sci.Symp.Conf.Rec. **2011**, 1226 (2011).
- [60] D. Y. Akimov *et al.*, Instrum.Exp.Tech. **55**, 423 (2012).
- [61] C. Duba *et al.*, J.Phys.Conf.Ser. **136**, 042077 (2008).
- [62] J. H. Davis, Phys.Rev.Lett. **113**, 081302 (2014), 1407.1052.
- [63] P. Barbeau, J. Collar, Y. Efremenko, and K. Scholberg, (2014), 1409.3185.